## Gauge theories on circles and the decoding of F-theory geometries



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Using results with Federico Bonetti, Stefan Hohenegger, Andreas Kapfer, Jan Keitel arXiv: 1112.1082, 1302.2918, 1305.1929

## Introduction / Motivation

## Effective actions of F-theory

- F-theory compactifications arguably provide the richest class of string theory effective actions
- original excitement about GUTs in F-theory - not realized in pert. strings

$$
\rightarrow \text { e.g. talk of Weigand }
$$

, numerous examples with exotic matter spectra $\rightarrow$ e.g. talks of Anderson, Klevers

- dualities to heterotic and Type I compactifications $\rightarrow$ e.g. talks of Mayrhofer, Lüst
' new 'exotic' theories, for example, 4D N=3 theories of [García-Etxebarria,Regalado]
$\rightarrow$ e.g. talk Iñaki
- Two questions arise:
- How can we infer reliably information about the 2D / 4D / 6D effective actions of F-theory?
- Is there a classification in sight? What do F-theory geometries classify?


## Some background material

- Well-known slogan: F-theory compactifications on elliptically-fibered Calabi-Yau threefolds yield 6D theories with
 minimal $N=(1,0)$ supersymmetry
pinching of two-torus indicates location of seven-branes
- brane and bulk physics encoded by singular complex geometry
- Six-dimensional theories are perfect to answer the above questions
- have a rich structure - there are many topologically distinct CY threefolds
- are strongly constraint by anomalies - fermions in 'all' $\mathrm{N}=1$ multiplets can contribute to anomalies


## Goals of this talk

- (1) Argue that the complete information in the F-theory geometries actually describe gauge/sugra theories on a circle.
- This is not unexpected since the M-theory to F-theory approach has been suggested already in [Vafa '96]. However, its importance and power might have been underappreciated. Currently this limit is the only reliable way to infer information about F-theory effective actions.
- (2) Show that F-theory geometries can teach us valuable lessons about circle-reduced theories.
- Example: How are anomalies of the higher-dimensional theories visible in the lower-dimensional effective theory?
- (3) Comment on possibility of classifying Calabi-Yau threefolds (with elliptic fibration) by using the insights from gauge theory.


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- (3) Comment on possibility of classifying Calabi-Yau threefolds (with elliptic fibration) by using the insights from gauge theory.

6D gauge theories / Sugra theories on circle $\Longleftrightarrow$
Geometry of resolved elliptically fibered Calabi-Yau threefolds
Our systematics is somewhat complementary to the approach and classifications of: [Morrison,Taylor etal.] [Heckman,Morrison,Rudelius,Vafa]

## Some comments on the M-theory to F-theory duality

## Approaching the problem via M-theory

- Viewing F-theory as an actually 12 dim. theory is problematic (e.g. torus volume unphysical, meaning of non-perturbative states...)
- F-theory effective actions via M-theory (as of now: definition of F-theory) Consider M-theory on space $T^{2} \times M_{9}$
$\tau$ is the complex structure modulus of the $T^{2}, \quad v$ volume of $T^{2}$


F-theory limit:
(1) A-cycle: if small than M-theory becomes Type IIA
(2) B-cycle: T-duality $\Rightarrow$ Type IIA becomes Type IIB, $\tau$ is indeed dilaton-axion
(3) grow extra dimension: send $v \rightarrow 0$ than T-dual B-cycle becomes large

## Consequences of M-theory to F-theory limit

- First step: approach M-theory via 11D supergravity on a smooth geometry
- resolution of singular Calabi-Yau geometry [almost everyone who has worked on $\Rightarrow$ classification of resolutions at each co-dimensions in base F-theory,...]
- co-dimension 1 in base
$\rightarrow$ non-Abelian gauge group
$\rightarrow$ simple roots
co-dimension 2 in base
$\rightarrow$ matter in representation $R$
$\rightarrow$ weights $w$ of $\boldsymbol{R}$

Abelian gauge group factors: $n_{U(1)}+1$ rational sections of the fibration zero section (assumed to exist throughout this talk)

- Second step: M-theory to F-theory limit shrinks fiber torus and resolutions and grows extra dimension (5D $\rightarrow 6 \mathrm{D}$ ) - keep track of M2-brane states
- e.g. M2-branes on two-torus fiber correspond to circle Kaluza-Klein states


## Massive states in five dimensions

- 11D Sugra on smooth geometry
$\Rightarrow$ 5D effective theory


Both types of theories need to be computed and then compared

- 5D gauge group in Coulomb branch: $U(1)^{\mathrm{rank} G} \times U(1)^{n_{U(1)}}$
- mass of 5D state at Kaluza-Klein level $n$ descending from 6D state in representation $\boldsymbol{R}$ of G with weights $w$ and $\mathrm{U}(1)$-charges $q_{m}$

$$
m=m_{\mathrm{CB}}+m_{\mathrm{KK}}=w_{I} \zeta^{I}+q_{m} \zeta^{m}+\frac{n}{r}
$$

Coulomb branch vevs (blow-up vevs) circle radius (torus fiber volume)

## Anomaly cancellation and circle compactifications

## M-theory on Calabi-Yau threefolds

- effective action has been studied long ago: 5D N=2 theory
- important to us are the Chern-Simons terms:

$$
-\frac{1}{48 \pi^{2}} \mathcal{K}_{A B C} \int_{\mathcal{M}_{5}} A^{A} \wedge F^{B} \wedge F^{C}-\frac{1}{384 \pi^{2}} c_{B} \int_{\mathcal{M}_{5}} A^{B} \wedge \operatorname{Tr}(R \wedge R)
$$

triple intersection numbers:
second Chern class: $A=1, \ldots, h^{1,1}\left(Y_{3}\right)$
$\mathcal{K}_{A B C}=\int_{Y_{3}} \omega_{A} \wedge \omega_{B} \wedge \omega_{C}$

$$
c_{A}=\int_{Y_{3}} \omega_{A} \wedge c_{2}\left(Y_{3}\right)
$$

- arise from 11D sugra action including terms up to 8 derivatives
- comparison to Chern-Simons terms obtained after circle reduction $\rightarrow$ classical and one-loop corrections required early works: [Morrison,Seiberg][Witten][Intriligator, Morrison,Seiberg]


## Chern-Simons terms on circle side

- Classical and one-loop Chern-Simons terms
- classical terms depend on $\Omega_{\alpha \beta}, b^{\alpha}, a^{\alpha}$ (6D tensor coupling, anomaly coefficients)
$\rightarrow \Omega_{\alpha \beta}, b^{\alpha}$, $a^{\alpha}$ fixed by geometry (intersection numbers base, location of branes, canonical class of base)
$\rightarrow$ there are many more Chern-Simons terms in M-theory
[Bonetti,TG '11]
- one-loop CS-terms in the effective theory induced by integrating out:
$\mathcal{K}_{A B C}=\sum_{\text {mass. states }} k_{\mathrm{r}} \cdot q_{A} q_{B} q_{C} \operatorname{sign}(m)$
massive spin 1 / 2 fermions massive spin $3 / 2$ fermions massive 'self-dual' tensors
[Bonetti,TG,Hohenegger '13]

$$
c_{A}=\sum_{\text {mass. states }} \kappa_{\mathrm{r}} \cdot q_{A} \operatorname{sign}(m)
$$

massive Kaluza-Klein states of all 6D fields that carry chirality and contribute to the 6D anomaly

## Jumping Chern-Simons terms

- match was still not possible for certain geometries

[TG,Kapfer,Keitel '13] regularization of infinite sum over KK modes in one-loop CS terms gets modified
- While one-loop CS terms are independent of the precise numerical value of the mass of a state, they do depend on
(1) sign of CB mass $\operatorname{sign}\left(m_{\mathrm{CB}}\right)$
(2) hierarchy of $m_{\mathrm{KK}}^{n}=n / r$ and $m_{\mathrm{CB}}$
$\rightarrow$ associate an integer label $\ell_{w, q}$ to each massive state: $m_{\mathrm{CB}}$ is between mass of $\ell_{w, q}$ and $\ell_{w, q}+1$ Kaluza-Klein state


## Extending box graphs

- Box graphs have been introduced to systematically classify the signinformation and realized Coulomb branch phases
[Hayashi, Lawrie, Morrison, Schäfer-Nameki][Braun,Schäfer-Nameki]...
- Example: anti-symmetric representation 10 of $\operatorname{SU}(5)$



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What are the rules for such diagrams? What are the inherent symmetries?

## Large gauge transformations

- large gauge transformations precisely induce certain types of jumps
, 6D gauge transformation: $\quad \hat{\Lambda}^{\mathcal{I}}(x, y)=\left\{\begin{array}{cc}-k^{I} y & \hat{\Lambda}^{m}(x, y)=-k^{m} y \\ 0\end{array}\right.$
, shift of Coulomb branch vevs: $\quad \tilde{\zeta}^{I}=\zeta^{I}+\frac{k^{I}}{r}, \quad \tilde{\zeta}^{m}=\zeta^{m}+\frac{k^{m}}{r}$,
mix of 5D vectors:

$$
\begin{aligned}
\tilde{A}^{I} & =A^{I}-k^{I} A^{0} \\
\tilde{A}^{m} & =A^{m}-k^{m} A^{0}
\end{aligned}
$$

mixing-in the Kaluza-Klein vector $A^{0}$ rearrangement of whole Kaluza-Klein tower

- large gauge transformations act on non-trivially on $\ell_{w, q} \Rightarrow$ Symmetry?
- Yes! But only if one takes into account 6D anomalies!

One-loop corrected effective 5D theories differing by the above transformation are identified if and only if 6D anomalies are cancelled.

## Discovering the 6D anomaly conditions

- Recall the form of the 6D anomaly cancellation conditions: e.g.
pure Abelian: $\quad-\left(b_{m n}^{\alpha} b_{p q}^{\beta}+b_{m p}^{\alpha} b_{n q}^{\beta}+b_{m q}^{\alpha} b_{n p}^{\beta}\right) \Omega_{\alpha \beta}=\sum_{q} F_{1 / 2}(q) q_{m} q_{n} q_{p} q_{q}$
pure non-Abelian:

$$
-3 \frac{b^{\alpha}}{\lambda(G)} \frac{b^{\beta}}{\lambda(G)} \Omega_{\alpha \beta}=\sum_{\boldsymbol{R}} F_{1 / 2}(\boldsymbol{R}) C_{\boldsymbol{R}}
$$

- differ from one-loop Chern-Simons terms: e.g.
pure Abelian:

$$
\mathcal{K}_{m n p}=\frac{1}{2} \sum_{q} F_{1 / 2}(q) q_{m} q_{n} q_{p}\left(2 \ell_{q}+1\right) \operatorname{sign}\left(m_{\mathrm{CB}}^{q}\right)
$$

pure non-Abelian:

$$
\mathcal{K}_{I J K}=\frac{1}{2} \sum_{\boldsymbol{R}} F_{1 / 2}(\boldsymbol{R}) \sum_{w \in \boldsymbol{R}} w_{I} w_{J} w_{K}\left(2 \ell_{w}+1\right) \operatorname{sign}\left(m_{\mathrm{CB}}^{w}\right)
$$

## Discovering the 6D anomaly conditions

- We were able to show:

Properties of 1-loop CS-terms:
(1) additional $q_{r}, w_{I}$ from large gauge transformation acting on

$$
\begin{aligned}
& \delta_{r}\left(2 \ell_{q}+1\right) \operatorname{sign}\left(m_{\mathrm{CB}}^{q}\right) \longrightarrow q_{r} \\
& \delta_{L}\left(2 \ell_{w}+1\right) \operatorname{sign}\left(m_{\mathrm{CB}}^{w}\right) \longrightarrow w_{L}
\end{aligned}
$$

(2) identities involving four weights to obtain $C_{\boldsymbol{R}}$

6D anomalies:

$$
\begin{aligned}
& \sum_{q} F_{1 / 2}(q) q_{m} q_{n} q_{p} q_{q} \\
& \sum_{\boldsymbol{R}} F_{1 / 2}(\boldsymbol{R}) C_{\boldsymbol{R}}
\end{aligned}
$$

[TG,Kapfer '15] Change in classical and 1-loop ChernSimons terms cancel if and only if 6D anomalies are cancelled

## Arithmetic structures and anomaly cancellation

- A proof anomaly cancellation for F-theory geometries requires to show that the large gauge transformations are actually a geometric symmetry.
- for models with only Abelian gauge symmetries this is possible [TG,Kapfer '15] large gauge transformations $\cong$ Mordell-Weil group of rational sections

$$
M W\left(Y_{3}\right) \cong \mathbb{Z}^{n_{U(1)}} \oplus \mathbb{Z}_{k_{1}} \oplus \ldots \oplus \mathbb{Z}_{k_{n_{\text {tor }}}}
$$

key: in the M-theory to F-theory limit we are free to pick the 'zero-section' that specifies the Kaluza-Klein vector (any choice works fine!)

- define new arithmetic structures for non-Abelian large gauge transformations define arithmetic structure on geometries with multi-sections
$\rightarrow$ mathematical meanings?
[TG,Kapfer,Klevers '15]
- Higgsing might connect non-Abelian models to Abelian models
$\rightarrow$ arithmetic structures need to be compatible


## Comments on the classification of threefolds

## Characteristic topological data

- Wall's classification theorem:
- The homotopy types of Calabi-Yau threefolds are classified by the following numerical characteristics:

$$
\begin{aligned}
& \text { Hodge numbers: } \quad h^{1,1}\left(Y_{3}\right), h^{2,1}\left(Y_{3}\right) \\
& \text { Triple intersection numbers: } \mathcal{K}_{A B C} \\
& \text { Second Chern class: } \quad c_{A}
\end{aligned}
$$

- Spaces differing in these numerical data cannot be continuously deformed into each other without encountering singularities.
- Are these data systematically specified by the F-theory effective actions on an additional circle?


## Information from gauge theories on circles

- (1) Hodge numbers are determined by spectrum (in the following: only non-Abelian)

$$
\begin{aligned}
& h^{1,1}\left(Y_{3}\right)=T+2+\operatorname{rank}(G) \\
& h^{2,1}\left(Y_{3}\right)=H_{\text {neut }}-1
\end{aligned}
$$

- (2) Intersection numbers and second Chern class determined by:
, classical data: $\Omega_{\alpha \beta}, b^{\alpha}, a^{\alpha}$

$$
\mathcal{K}_{0 \alpha \beta}=\Omega_{\alpha \beta} \quad \mathcal{K}_{I J \alpha}=-\mathcal{C}_{I J} \Omega_{\alpha \beta} b^{\beta} \quad c_{\alpha}=-12 \Omega_{\alpha \beta} a^{\beta}
$$

- one-loop data: spectrum plus circle info (sign table, jump levels)

$\rightarrow$ extended box graphs as convenient object for classification?


## Intersection and Chern class data

- Intersection numbers and Chern classes (non-Abelian groups only)

$$
\begin{aligned}
\mathcal{K}_{000} & =\frac{1}{120}\left(2\left(T_{s d}-T_{a s d}\right)-F_{1 / 2}-5 F_{3 / 2}\right)+\frac{1}{4} \sum_{R} F_{1 / 2}(R) \sum_{w \in R} l_{w}^{2}\left(l_{w}+1\right)^{2}, \\
\mathcal{K}_{00 I} & =\frac{1}{6} \sum_{R} F_{1 / 2}(R) \sum_{w \in R} l_{w}\left(l_{w}+1\right)\left(2 l_{w}+1\right) w_{I} \operatorname{sign}\left(m_{\mathrm{CB}}^{w}\right), \\
\mathcal{K}_{0 I J} & =\frac{1}{12} \sum_{R} F_{1 / 2}(R) \sum_{w \in R}\left(1+6 l_{w}\left(l_{w}+1\right)\right) w_{I} w_{J}, \\
\mathcal{K}_{I J K} & =\frac{1}{2} \sum_{R} F_{1 / 2}(R) \sum_{w \in R}\left(2 l_{w}+1\right) w_{I} w_{J} w_{K} \operatorname{sign}\left(m_{\mathrm{CB}}^{w}\right), \\
c_{0} & =\frac{1}{6}\left(19 F_{3 / 2}-F_{1 / 2}-4\left(T_{s d}-T_{a s d}\right)\right)-\sum_{R} F_{1 / 2}(R) \sum_{w \in R} l_{w}\left(l_{w}+1\right), \\
c_{I} & =-\sum_{R} F_{1 / 2}(R) \sum_{w \in R}\left(2 l_{w}+1\right) w_{I} \operatorname{sign}\left(m_{\mathrm{CB}}^{w}\right),
\end{aligned}
$$

## Two classification problems

- (1) given a 6D N=(1,0) anomaly-free theory in tensor Coulomb branch $\rightarrow$ construct classifying topological data of Calabi-Yau threefolds
- Systematics: - start with gauge theories of non-Higgsable clusters $\begin{gathered}\text { [Morrison, } \\ \text { Taylor] }\end{gathered}$
- successive un-Higgsing of gauge groups in field theory
- generate tree of topological data
- Challenges: - moding out symmetries
- detecting additional geometric constraints
- extending analysis to non-resolvable geometries
$\Rightarrow$ String Universality: Is F-theory a theory of F-rything in 6D?
- (2) given an set of Calabi-Yau geometries, as provided e.g. by the Kreuzer-Skarke list
$\rightarrow$ classify $6 \mathrm{D} \mathrm{N}=(1,0)$ theories associated to elliptic fibrations


## Conclusions

- The map between the Calabi-Yau threefold geometry and a 6D sugra theory on a circle has been established at classical and one-loop level.
- Anomalies can be inferred from Chern-Simons terms, if the latter are probed by large gauge transformations.
- Large gauge transformations map to arithmetic structures on ell. fibrations.
for Abelian gauge groups one encounters the Mordell-Weil group
$\rightarrow$ general proof of anomaly cancellation for viable F-theory geometries
new arithmetic structures for non-Abelian groups and geometries with multi-sections.
, Extension to 3D / 4D has been worked out for chiral spectrum induced by fluxes
- F-theory provides a novel way to approach classification problems.
- topological data of resolved Calabi-Yau threefolds are robust and corse information directly related to corse information about gauge theories

Are there mathematical restrictions on viable Hodge numbers, intersection numbers Chern classes? $\rightarrow$ ruling out 6D gauge theories coupled to gravity?

## Thank you for your attention!

