# Extremal transitions of Calabi-Yau fourfolds in M-theory 

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F-Theory at 20: Dave Day February 25, 2016<br>H. Jockers, SK, D.R. Morrison, M.R. Plesser<br>arXiv:1602.xxxxx

## Outline

(1) Extremal Transitions in M-Theory

2 Field Theory Analysis
(3) Toric Case

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## (2) Field Theory Analysis

(3) Toric Case

- $X$ CY3, $M[X]$
- 3-form gauge field $C_{3}, G=d C_{3}$
- $G$ is quantized:

- Tadpole cancellation

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M=\frac{\chi(X)}{24}-\frac{1}{2} \int_{X} \frac{G}{2 \pi} \wedge \frac{G}{2 \pi} \in \mathbb{Z}
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## Extremal Transitions in M-Theory

- From singular $X_{0}$, can blow up to $X^{\sharp}$ or deform to $X^{b}$

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\begin{array}{ccc} 
& & X^{\sharp} \\
& & \downarrow \\
X^{b} \rightsquigarrow & X_{0}
\end{array}
$$

- In M-theory, need $G^{\sharp}$ and $G^{b}$
- $M^{\sharp}$ and $M^{b} \mathrm{M} 2$-branes for tadpole cancellation
- Look for transitions for which the M2 branes are spectators-M2-branes kept away from $S=\operatorname{Sing}\left(X_{0}\right)$



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& & \\
& \\
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$$
\frac{\chi\left(X^{b}\right)}{24}-\frac{\chi\left(X^{\sharp}\right)}{24}=\frac{1}{2} \int_{X} \frac{G^{b}}{2 \pi} \wedge \frac{G^{b}}{2 \pi}-\frac{1}{2} \int_{X} \frac{G^{\sharp}}{2 \pi} \wedge \frac{G^{\sharp}}{2 \pi}
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- Extremal transition where $S \subset X_{0}$ is a smooth surface of transverse $A_{n-1}$ singularities
- Local equation $x y=z^{n}, z \in K_{S}$
- If $c_{2}\left(X^{\sharp}\right)$ is even, then $G^{\sharp}=0$ satisfies quantization and tadpole cancellation for suitable $M^{\sharp}$
- The main result in this talk: given any toric hypersurface $X_{0}$ such that $c_{2}\left(X^{\sharp}\right)$ is the restriction of an even class on the toric variety, we can always find a transition in M theory from $\left(X^{\sharp}, G^{\sharp}=0\right)$ to an $\left(X^{j}, G^{b}\right)$ which satisfies quantization and tadpole cancellation.
- Furthermore, the geometric moduli of $\left(X^{b}, G^{b}\right)$ perfectly match the predictions of the low energy theory
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## Constraints from M theory

- Constraints from M theory on transition $\left(X^{\sharp}, 0\right) \rightarrow\left(X^{b}, G^{b}\right)$ :
- Quantization:

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using

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\chi\left(X^{b}\right)-\chi\left(X^{\sharp}\right)=(n+1) n(n-1) K_{S}^{2}
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\frac{1}{2} \int_{X} \frac{G^{b}}{2 \pi} \wedge \frac{G^{b}}{2 \pi}=\frac{(n+1) n(n-1) K_{S}^{2}}{24}
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(2) Field Theory Analysis
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## Low Energy Theory

- Low energy 3D dynamics is given by a twisted dimensional reduction of $N=17 \mathrm{D} S U(n)$ SYM on $S$
- $S U(n)$ gauge theory with $p_{g}+q$ adjoint chirals

$$
p_{g}=h^{2,0}(S)=h^{0}\left(S, K_{S}\right), \quad q=h^{1,0}(S)
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- Coulomb branch of dimension $(n-1)\left(p_{g}+q+1\right)$, parametrized by vevs of gauge bosons in Cartan

and the vevs of the $n-1 U(1)^{n-1}$-neutral scalars in each
of the $p_{g}+q$ chirals
- Residual $S_{n}=W\left(A_{n-1}\right)$ action on Coulomb branch
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## Checks against gauge theory predictions

- $(n-1)\left(p_{g}+1\right)$ Columb branch moduli $(q=0)$ : blowup modes in Kähler moduli and

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x y+\prod_{i=1}^{n}\left(z+\eta_{i}\right), \quad \eta_{i} \in H^{0}\left(S, K_{S}\right), \sum \eta_{i}=0
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## Toric CY 4folds with $A_{n-1}$ singularities

- $N, M$ dual lattices of rank 5
- $\left(\Delta, \Delta^{\circ}\right)$ polar 5-dimensional reflexive polytopes
- $\Delta \subset M_{\mathbb{R}}, \Delta^{\circ} \subset N_{\mathbb{R}}$
- 0 unique interior point of $\Delta \cap M$ and of $\Delta^{\circ} \cap N$
- Fan $\Sigma \sharp$ of toric variety $\mathbb{P}_{\Delta}$ : cones over the faces of $\Delta$
- Highly singular, so we choose a maximal projective crepant subdivision of that fan
- $A_{n-1}$ case: $\Delta^{\circ}$ has a one-dimensional edge $\Gamma$ containing $n-1$ interior lattice points
- $X^{\sharp} \subset X_{\Sigma \sharp}$ anticanonical hypersurface - CY 4fold
- Removing the interior lattice points of $\Gamma$ blows down $X_{\Sigma^{\sharp}}$ to $X_{\Sigma_{0}}$, and $X^{\sharp}$ to $X_{0}$


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## Example

- Describe $\mathbf{P}(1,1,2,2,2,2)$ (cf. Ronen's talk) torically
- Take edges of $\Sigma_{0}$ as rows of

- Have unique interior point $(0,-1,-1,-1,-1)$ of edge joining first two vertices
- SU(2) gauge theory
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- Have unique interior point $(0,-1,-1,-1,-1)$ of edge joining first two vertices
- SU(2) gauge theory
- Label vertices of $\Sigma_{0}$ as $v_{1}, \ldots, v_{k} ; v_{1}, v_{2}$ endpoints of $\Gamma$
- For each $v_{j}$ we have homogeneous coordinate $x_{j}$ and divisor $D_{j} \subset X_{\Sigma_{0}}$ given by $x_{j}=0$
- Interior points $v_{0}, v_{-1}, \ldots, v_{2-n}$ in order
- For the vertices $v_{2-n}, \ldots, v_{k}$ we similarly have toric divisors $D_{j}^{\sharp} \subset X^{\Sigma^{\sharp}}$
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- For the vertices $v_{2-n}, \ldots, v_{k}$ we similarly have toric divisors $D_{j}^{\sharp} \subset X^{\Sigma^{\sharp}}$
- Embed $\iota: X_{\Sigma} \hookrightarrow X_{\Sigma^{b}}$ for suitable fan $\Sigma^{b}$ : $\left(x_{1}, x_{2}, x_{3}, \ldots, x_{k}\right) \mapsto\left(x_{1}^{n}, x_{2}^{n}, x_{1} x_{2}, x_{3}, \ldots, x_{k}\right)=:$ $\left(y_{0}, \ldots, y_{k}\right)$
- The $A_{n-1}$ is visible from $q_{0}(y):=y_{0} y_{1}-y_{2}^{n}=0$
- Convenient choice for $\Sigma^{b}$ : choose $m_{\Gamma} \in M$

$$
\left\langle m_{\Gamma}, v_{1}\right\rangle=n-1,\left\langle m_{\Gamma}, v_{2}\right\rangle=-1
$$

- Take as edges of $\Sigma^{b}$, with appropriate higher-dimensional cones

$$
\begin{aligned}
& w_{0}=\left(\frac{v_{1}-v_{2}}{n},-(n-1)\right), \\
& w_{1}=(0,1) \\
& w_{2}=\left(v_{2}, 0\right), \\
& w_{i}=\left(v_{i},-n\left\langle m_{\Gamma}, v_{i}\right\rangle\right), i \geq 3
\end{aligned}
$$

- Embed $\iota: X_{\Sigma} \hookrightarrow X_{\Sigma^{b}}$ for suitable fan $\Sigma^{b}$ : $\left(x_{1}, x_{2}, x_{3}, \ldots, x_{k}\right) \mapsto\left(x_{1}^{n}, x_{2}^{n}, x_{1} x_{2}, x_{3}, \ldots, x_{k}\right)=:$ $\left(y_{0}, \ldots, y_{k}\right)$
- The $A_{n-1}$ is visible from $q_{0}(y):=y_{0} y_{1}-y_{2}^{n}=0$


## - Convenient choice for $\sum^{\bullet}$ : choose $m_{\Gamma} \in M$

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\left\langle m_{\Gamma}, v_{1}\right\rangle=n-1,\left\langle m_{\Gamma}, v_{2}\right\rangle=-1
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$$
\begin{aligned}
w_{0} & =\left(\frac{v_{1}-v_{2}}{n},-(n-1)\right) \\
w_{1} & =(0,1) \\
w_{2} & =\left(v_{2}, 0\right) \\
w_{i} & =\left(v_{i},-n\left\langle m_{\Gamma}, v_{i}\right\rangle\right), i \geq 3
\end{aligned}
$$

- Have toric divisors $D_{j}^{b} \subset X_{\Sigma^{b}}$ associated with edges $w_{j}$
- $D_{0}^{3}+D_{1}^{p} \sim N D_{2}^{p}$

- $\iota^{*}\left(D^{\prime}\right)=D_{1}+\ldots+D_{k}, D^{\prime}:=D_{2}+D_{3}+\ldots+D_{k}$
- Choose section $g(y)$ of $D^{\prime}$ which pulls back by $\iota$ to an equation for $X_{0}$
- $\iota\left(X_{0}\right)$ is the complete intersection of $q(y)$ and $g(y)$
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## Example

- Return to $\mathbf{P}(1,1,2,2,2,2)$. Take $m_{\Gamma}=(1,0,0,0,0)$
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\left(\begin{array}{cccccc}
-1 & -1 & -1 & -1 & -1 & -1 \\
0 & 0 & 0 & 0 & 0 & 1 \\
1 & 0 & 0 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 & 0 & 0 \\
0 & 0 & 1 & 0 & 0 & 0 \\
0 & 0 & 0 & 1 & 0 & 0 \\
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- $X_{0}$ is a (2.5) complete intersection in $P^{6}$


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- Fan for $\mathbf{P}^{6}$
- $X_{0}$ is a $(2,5)$ complete intersection in $\mathbf{P}^{6}$
- $S \subset X_{\Sigma^{b}}$ is the complete intersection of $y_{0}, y_{1}, y_{2}, g(y)$
- $K_{S}=\left.K_{X_{\Sigma} b}\left(D_{0}^{b}+D_{1}^{p}+D_{2}^{b}+D^{\prime}\right)\right|_{S}=\mathcal{O}_{S}\left(D_{2}\right)$
- Basis for sections of $K_{S}$ correspond to interior lattice points of dual face $\Gamma^{\circ}$
- $p_{g}$ is the number of these points.
- Relabel coordinates so these correspond to $y_{3}, \ldots, y_{p_{g}+2}$
- $\tilde{a}=y_{0} y_{1}-y_{2}^{n}-\sum_{j=0}^{n-2} h_{n-j}\left(y_{3}, \ldots, y_{p_{g}+2}\right) y_{2}^{j}$, deg $h_{j}=j$
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- Will exhibit $G^{b}$ with the following smoothings of $X_{0}$ as flat directions

$$
\tilde{q}=y_{0} y_{1}-\operatorname{det}\left(y_{2} I_{n}+M\left(y_{3}, \ldots, y_{p_{g}+2}\right)\right)
$$

with $M(y)$ a traceless $n \times n$ matrix of linear forms

- $\left(p_{g}-1\right)\left(n^{2}-1\right)$ moduli for $\tilde{q}:$
- $p_{g}\left(n^{2}-1\right)$ for entries of $M(y)$
- $n^{2}-1$ similarity transformations of $M(y)$ leave $\tilde{q}$ unchanged
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## - Specialize to $n=2$ for simplicity



- $T_{1} \subset X_{\Sigma^{b}}$ defined by $y_{0}=y_{2}+\ell_{11}(y)=\ell_{12}(y)=g(y)=0$
- By construction, $T_{1} \subset X^{\prime}$ since the first row of $y_{2} l_{2}+M(y)$ is $\left(y_{2}+\ell_{11}(y), \ell_{12}(y)\right)$
- Similarly $T_{2} \subset X^{b}$ defined by
$y_{0}=y_{2}+\ell_{11}(y)=\ell_{21}(y)=g(y)=0$

$$
\frac{G^{b}}{2 \pi}:=\frac{1}{2}\left(T_{1}-T_{2}\right) \in H^{4}\left(X^{b}, \mathbf{R}\right)
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with $2 \frac{G^{b}}{2 \pi} \in H^{4}\left(X^{b}, Z\right)$

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- $G^{b}$ is an algebraic class, hence of type $(2,2)$
- Let $F \subset X_{\Sigma}$ be the hypersurface $g(y)=0$ - $T_{1}, T_{2}$ complete intersections in $F$ of same degrees - Therefore image of $G^{b}$ in $H^{6}(F)$ vanishes - $G^{b}$ primitive
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- Let $F \subset X_{\Sigma^{\triangleright}}$ be the hypersurface $g(y)=0$
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- Therefore image of $G^{\prime}$ in $H^{6}(F)$ vanishes
- G' primitive
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- $G^{b}$ primitive
- Tadpole: computing in $X_{\Sigma^{b}}$

$$
\int_{X^{b}} \frac{G^{b}}{2 \pi} \wedge \frac{G^{b}}{2 \pi}=\frac{(n+1) n(n-1)}{12} \int_{X_{\Sigma}^{b}} D_{0}^{b} D_{1}^{b}\left(D_{2}^{b}\right)^{3} D^{\prime}
$$

- Recall: $S$ is the complete of divisors in the classes $D_{0}^{b}, D_{1}^{b}, D_{2}^{p}, D^{\prime}$
- $K_{S}$ is the restriction of $D_{2}$
- Integral on the right is just $K_{S}^{2}$
- Divide by two to verify tadpole cancellation


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- Recall: $S$ is the complete of divisors in the classes
$D_{0}^{b}, D_{1}^{b}, D_{2}^{b}, D^{\prime}$
- $K_{S}$ is the restriction of $D_{2}$
- Integral on the right is just $K_{S}^{2}$
- Divide by two to verify tadpole cancellation


## Quantization

$$
0 \rightarrow \mathcal{O}_{X_{\Sigma^{\sharp}}^{k+n-6}}^{k+6} \rightarrow \bigoplus_{i=2-n}^{k} \mathcal{O}_{x_{\Sigma^{\sharp}}}\left(D_{i}^{\sharp}\right) \rightarrow T_{X_{\Sigma^{\sharp}}} \rightarrow 0
$$

- Leading to $c_{2}\left(X^{\sharp}\right)=\sum_{i<j} D_{i}^{\sharp} D_{j}^{\sharp}$
- Similarly $c_{2}\left(X^{b}\right)=\sum_{i<j} D_{i}^{b} D_{j}^{b}$

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\frac{G^{b}}{2 \pi}=\frac{1}{2}\left(T_{1}-T_{2}\right) \equiv \frac{1}{2}\left(T_{1}+T_{2}\right)(\bmod \mathbf{Z})
$$

- But $T_{1}+T_{2}$ is the complete intersection $y_{2}+\ell_{11}(y)=\tilde{q}(y)$ in $F$, cohomology class $\left(D_{2}^{b}\right)^{2}$
- We can replace $G^{j} /(2 \pi)$ with $\left(D_{2}^{j}\right)^{2} / 2$
- Everything is now explicitly computable and we verify quantization

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- The Weyl group action permutes the ordering of the rows and columns of $M(y), T_{1} \leftrightarrow T_{2}$
- Coulomb branch: $M(y)=\operatorname{diag}(\eta(y),-\eta(y))$, permuting ordering of rows and columns
- Agrees with Weyl group action on Coulomb branch
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## HAPPY BIRTHDAY DAVE!


[^0]:    - Fan for $\mathbf{P}^{6}$
    - $X_{0}$ is a $(2,5)$ complete intersection in $\mathrm{P}^{6}$

