Extremal transitions of Calabi-Yau fourfolds in M-theory

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• X CY3, *M*[X]

- 3-form gauge field C_3 , $G = dC_3$
- *G* is quantized:

$$rac{G}{2\pi}-rac{c_2(X)}{2}\in H^4(M,{f Z})$$

• Tadpole cancellation

$$M=rac{\chi(X)}{24}-rac{1}{2}\int_Xrac{G}{2\pi}\wedgerac{G}{2\pi}\in{\sf Z},$$

number of M2 branes

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Superpotentials

• G induces superpotential

$$W = \int_X \Omega \wedge rac{G}{2\pi}$$

and twisted superpotential

$$ilde{W} = \int_X J \wedge J \wedge rac{G}{2\pi}$$

• Flat directions: $G \in H^{2,2}(X)$, *G* primitive

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• From singular X_0 , can blow up to X^{\sharp} or deform to X^{\flat}



- In M-theory, need G^{\sharp} and G^{\flat}
- M^{\sharp} and M^{\flat} M2-branes for tadpole cancellation
- Look for transitions for which the M2 branes are spectators—M2-branes kept away from S = Sing(X₀)
- $M^{\sharp} = M^{\flat}$, or

$$\frac{\chi(X^{\flat})}{24} - \frac{\chi(X^{\sharp})}{24} = \frac{1}{2} \int_X \frac{G^{\flat}}{2\pi} \wedge \frac{G^{\flat}}{2\pi} - \frac{1}{2} \int_X \frac{G^{\sharp}}{2\pi} \wedge \frac{G^{\sharp}}{2\pi}$$

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- Extremal transition where S ⊂ X₀ is a smooth surface of transverse A_{n-1} singularities
- Local equation $xy = z^n$, $z \in K_S$
- If c₂(X[♯]) is even, then G[♯] = 0 satisfies quantization and tadpole cancellation for suitable M[♯]
- The **main result** in this talk: given any toric hypersurface X_0 such that $c_2(X^{\sharp})$ is the restriction of an even class on the toric variety, we can always find a transition in M theory from $(X^{\sharp}, G^{\sharp} = 0)$ to an (X^{\flat}, G^{\flat}) which satisfies quantization and tadpole cancellation.
- Furthermore, the geometric moduli of (X^b, G^b) perfectly match the predictions of the low energy theory

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Constraints from M theory

Constraints from M theory on transition (X[♯], 0) → (X[♭], G[♭]):
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Tadpole cancellation

$$\frac{1}{2}\int_X \frac{G^\flat}{2\pi} \wedge \frac{G^\flat}{2\pi} = \frac{(n+1)n(n-1)K_S^2}{24}$$

using

$$\chi(X^{\flat}) - \chi(X^{\sharp}) = (n+1)n(n-1)K_S^2$$

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Sheldon Katz Extremal transitions of Calabi-Yau fourfolds in M-theory

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• Low energy 3D dynamics is given by a twisted dimensional reduction of N = 1 7D SU(n) SYM on S

• SU(n) gauge theory with $p_g + q$ adjoint chirals

$$p_g = h^{2,0}(S) = h^0(S, K_S), \qquad q = h^{1,0}(S)$$

• Coulomb branch of dimension $(n-1)(p_g + q + 1)$, parametrized by vevs of gauge bosons in Cartan

$$(\phi_1,\ldots,\phi_n), \qquad \sum \phi_i = 0$$

and the vevs of the $n - 1 U(1)^{n-1}$ -neutral scalars in each of the $p_g + q$ chirals

- Residual $S_n = W(A_{n-1})$ action on Coulomb branch
- Higgs branch dimension $(n^2 1)(p_g + q 1)$

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(n − 1)(p_g + 1) Columb branch moduli (q = 0): blowup modes in Kähler moduli and

$$xy + \prod_{i=1}^{n} (z + \eta_i), \qquad \eta_i \in H^0(\mathcal{S}, \mathcal{K}_{\mathcal{S}}), \ \sum \eta_i = 0$$

- $(n^2 1)(p_g 1)$ flat directions for G^{\flat}
- Have *S_n* action on Coulomb branch
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Sheldon Katz Extremal transitions of Calabi-Yau fourfolds in M-theory

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Toric CY 4 folds with A_{n-1} singularities

- N, M dual lattices of rank 5
- (Δ, Δ°) polar 5-dimensional reflexive polytopes
- $\Delta \subset M_{\mathbb{R}}, \Delta^{\circ} \subset N_{\mathbb{R}}$
- 0 unique interior point of $\Delta \cap M$ and of $\Delta^{\circ} \cap N$
- Fan Σ^{\sharp} of toric variety \mathbb{P}_{Δ} : cones over the faces of Δ°
- Highly singular, so we choose a maximal projective crepant subdivision of that fan
- A_{n-1} case: Δ° has a one-dimensional edge Γ containing n − 1 interior lattice points
- $X^{\sharp} \subset X_{\Sigma^{\sharp}}$ anticanonical hypersurface CY 4fold
- Removing the interior lattice points of Γ blows down $X_{\Sigma^{\sharp}}$ to $X_{\Sigma_0},$ and X^{\sharp} to X_0

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Describe P(1, 1, 2, 2, 2, 2) (cf. Ronen's talk) torically
 Take edges of Σ₀ as rows of



- Have unique interior point (0, -1, -1, -1, -1) of edge joining first two vertices
- *SU*(2) gauge theory

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- Describe P(1, 1, 2, 2, 2, 2) (cf. Ronen's talk) torically
- Take edges of Σ₀ as rows of

- Have unique interior point (0, -1, -1, -1, -1) of edge joining first two vertices
- SU(2) gauge theory

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Label vertices of Σ₀ as v₁,..., v_k; v₁, v₂ endpoints of Γ

- For each v_j we have homogeneous coordinate x_j and divisor D_j ⊂ X_{Σ₀} given by x_j = 0
- Interior points $v_0, v_{-1}, \ldots, v_{2-n}$ in order
- For the vertices v_{2-n}, \ldots, v_k we similarly have toric divisors $D_j^{\sharp} \subset X^{\Sigma^{\sharp}}$

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- Embed $\iota : X_{\Sigma} \hookrightarrow X_{\Sigma^{\flat}}$ for suitable fan Σ^{\flat} : $(x_1, x_2, x_3, \dots, x_k) \mapsto (x_1^n, x_2^n, x_1 x_2, x_3, \dots, x_k) =:$ (y_0, \dots, y_k)
- The A_{n-1} is visible from $q_0(y) := y_0 y_1 y_2^n = 0$
- Convenient choice for Σ^{\flat} : choose $m_{\Gamma} \in M$

$$\langle m_{\Gamma}, v_1 \rangle = n - 1, \ \langle m_{\Gamma}, v_2 \rangle = -1$$

$$w_{0} = \left(\frac{v_{1} - v_{2}}{n}, -(n-1)\right) ,$$

$$w_{1} = (0, 1) ,$$

$$w_{2} = (v_{2}, 0) ,$$

$$w_{i} = (v_{i}, -n\langle m_{\Gamma}, v_{i} \rangle) , i \geq 3 .$$

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- Have toric divisors D^b_i ⊂ X_{Σ^b} associated with edges w_j
- $D_0^{\flat} + D_1^{\flat} \sim ND_2^{\flat}$
- $\iota^*(D_2^\flat) = D_1 + D_2$
- $\iota^*(D') = D_1 + \ldots + D_k, \ D' := D_2 + D_3 + \ldots + D_k$
- Choose section g(y) of D' which pulls back by ι to an equation for X_0
- $\iota(X_0)$ is the complete intersection of q(y) and g(y)
- X^{\flat} (w/o flux constraint) obtained by smoothing $\tilde{q}(y)$ of q(y)



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Return to P(1, 1, 2, 2, 2, 2). Take m_Γ = (1, 0, 0, 0, 0)
 Edges of Σ^b are rows of



- Fan for **P**⁶
- X_0 is a (2,5) complete intersection in \mathbf{P}^6

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$$\left(\begin{array}{cccccccc} -1 & -1 & -1 & -1 & -1 & -1 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \end{array}\right)$$

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- Fan for P⁶
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K_S

• $S \subset X_{\Sigma^{\flat}}$ is the complete intersection of $y_0, y_1, y_2, g(y)$

- $K_S = K_{X_{\Sigma^{\flat}}}(D_0^{\flat} + D_1^{\flat} + D_2^{\flat} + D')|_S = \mathcal{O}_S(D_2)$
- Basis for sections of K_S correspond to interior lattice points of dual face Γ°
- p_g is the number of these points.
- Relabel coordinates so these correspond to y_3, \ldots, y_{p_q+2}

•
$$\tilde{q} = y_0 y_1 - y_2^n - \sum_{j=0}^{n-2} h_{n-j}(y_3, \dots, y_{p_g+2}) y_2^j$$
, deg $h_j = j$

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Will exhibit G^b with the following smoothings of X₀ as flat directions

$$ilde{q} = y_0 y_1 - \det \left(y_2 I_n + M(y_3, \dots, y_{
ho_g+2})
ight)$$

with M(y) a traceless $n \times n$ matrix of linear forms

- $(p_g 1)(n^2 1)$ moduli for \tilde{q} :
- $p_g(n^2 1)$ for entries of M(y)
- n² 1 similarity transformations of M(y) leave q
 unchanged
- Matches Higgs branch moduli perfectly!

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Flat Directions

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• Specialize to *n* = 2 for simplicity

$$M(y) = \begin{pmatrix} \ell_{11}(y) & \ell_{12}(y) \\ \ell_{21}(y) & -\ell_{11}(y) \end{pmatrix}$$

- $T_1 \subset X_{\Sigma^{\flat}}$ defined by $y_0 = y_2 + \ell_{11}(y) = \ell_{12}(y) = g(y) = 0$
- By construction, $T_1 \subset X^{\flat}$ since the first row of $y_2 I_2 + M(y)$ is $(y_2 + \ell_{11}(y), \ell_{12}(y))$
- Similarly T₂ ⊂ X^b defined by y₀ = y₂ + ℓ₁₁(y) = ℓ₂₁(y) = g(y) = 0

$$rac{G^{\flat}}{2\pi}:=rac{1}{2}\left(T_{1}-T_{2}
ight)\in H^{4}(X^{\flat},\mathbf{R})$$

with $2rac{G^{\flat}}{2\pi}\in H^4(X^{\flat},\mathbf{Z})$



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• G^{\flat} is an algebraic class, hence of type (2, 2)

- Let $F \subset X_{\Sigma^{\flat}}$ be the hypersurface g(y) = 0
- T₁, T₂ complete intersections in F of same degrees
- Therefore image of G^{\flat} in $H^{6}(F)$ vanishes
- G[♭] primitive

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Tadpole

Tadpole: computing in X_{Σ^b}

$$\int_{X^{\flat}} \frac{G^{\flat}}{2\pi} \wedge \frac{G^{\flat}}{2\pi} = \frac{(n+1)n(n-1)}{12} \int_{X^{\flat}_{\Sigma}} D^{\flat}_0 D^{\flat}_1 (D^{\flat}_2)^3 D'$$

- Recall: S is the complete of divisors in the classes D₀^b, D₁^b, D₂^b, D'
- *K_S* is the restriction of *D*₂
- Integral on the right is just K_S^2
- Divide by two to verify tadpole cancellation

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- Leading to $c_2(X^{\sharp}) = \sum_{i < j} D_i^{\sharp} D_j^{\sharp}$
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- The Weyl group action permutes the ordering of the rows and columns of *M*(*y*), *T*₁ ↔ *T*₂
- $G^{\flat} \mapsto -G^{\flat}$
- Coulomb branch: M(y) = diag(η(y), -η(y)), permuting ordering of rows and columns
- $\eta \mapsto -\eta$
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HAPPY BIRTHDAY DAVE!

Sheldon Katz Extremal transitions of Calabi-Yau fourfolds in M-theory

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