

Extremal Transitions with Enhanced Gauge Symmetry in M-theory on Calabi–Yau Fourfolds

M. Ronen Plesser

F Theory at 20 – Dave Day
Caltech, 2/25/2016

H. Jockers, S. Katz, D.R. Morrison, MRP, arXiv:1602.xxxx

Outline

- 1 Motivation
- 2 M-Theory Compactifications, Fluxes, and Superpotentials
- 3 Codimension-two Degenerations and Gauge Theory Predictions
- 4 The Geometric Transition
- 5 An Example

Singular Geometries as Probes of Physics

- Singular compactifications provide insight into non-perturbative effects

Singular Geometries as Probes of Physics

- Singular compactifications provide insight into non-perturbative effects
- Near a singularity solitons wrapping shrinking cycles become light and participate in low-energy dynamics **Witten, Strominger**

Singular Geometries as Probes of Physics

- Singular compactifications provide insight into non-perturbative effects
- Near a singularity solitons wrapping shrinking cycles become light and participate in low-energy dynamics **Witten, Strominger**
- At low energy and near the singularity, bulk dynamics decouple leaving a field theory of local modes

Singular Geometries as Probes of Physics

- Singular compactifications provide insight into non-perturbative effects
- Near a singularity solitons wrapping shrinking cycles become light and participate in low-energy dynamics **Witten, Strominger**
- At low energy and near the singularity, bulk dynamics decouple leaving a field theory of local modes
- Geometric Moduli Space \leftrightarrow Field Theory Moduli Space

Singular Geometries as Probes of Physics

- Singular compactifications provide insight into non-perturbative effects
- Near a singularity solitons wrapping shrinking cycles become light and participate in low-energy dynamics **Witten, Strominger**
- At low energy and near the singularity, bulk dynamics decouple leaving a field theory of local modes
- Geometric Moduli Space \leftrightarrow Field Theory Moduli Space

Extremal Transitions

- As Kähler class of X^\sharp degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^b

Extremal Transitions

- As Kähler class of X^\sharp degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^b
- **Reid:** All Calabi–Yau spaces can be connected by extremal transitions

Extremal Transitions

- As Kähler class of X^\sharp degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^b
- **Reid**: All* Calabi–Yau spaces can be connected by extremal transitions

Extremal Transitions

- As Kähler class of X^\sharp degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^b
- **Reid**: All* Calabi–Yau spaces can be connected by extremal transitions
- Physically, transition is implemented by soliton condensation **Greene, Morrison, Strominger**

Extremal Transitions

- As Kähler class of X^\sharp degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^b
- Reid: All* Calabi–Yau spaces can be connected by extremal transitions
- Physically, transition is implemented by soliton condensation Greene, Morrison, Strominger
- Transitions in codimension-two involve enhanced gauge symmetry in type-II strings Aspinwall,Louis; Klemm, Mayr; Katz, Morrison,MRP

Extremal Transitions

- As Kähler class of X^\sharp degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^b
- Reid: All* Calabi–Yau spaces can be connected by extremal transitions
- Physically, transition is implemented by soliton condensation Greene, Morrison, Strominger
- Transitions in codimension-two involve enhanced gauge symmetry in type-II strings Aspinwall, Louis; Klemm, Mayr; Katz, Morrison, MRP
- M-theory on CY threefold singularities can teach us about five-dimensional QFT Seiberg, Morrison

Extremal Transitions

- As Kähler class of X^{\sharp} degenerates, singularity acquired can be smoothed by deformation to a topologically distinct Calabi–Yau space X^{\flat}
- **Reid**: All* Calabi–Yau spaces can be connected by extremal transitions
- Physically, transition is implemented by soliton condensation **Greene, Morrison, Strominger**
- Transitions in codimension-two involve enhanced gauge symmetry in type-II strings **Aspinwall, Louis; Klemm, Mayr; Katz, Morrison, MRP**
- M-theory on CY threefold singularities can teach us about five-dimensional QFT **Seiberg, Morrison**
- In M-theory on CY fourfolds fourform flux G plays an essential role, can use transitions to better understand it **Intriligator, Jockers, Katz, Mayr, Morrison, MRP**

M-theory Compactified on CY Fourfold X

- Spectrum ($\mathcal{N} = 2$):
 - $U(1)^{h^{1,1}(X)}$ gauge symmetry - scalars parameterize Kähler class
 $J = \sum x_i e^i$
 - $h^{3,1}(X)$ neutral chiral multiplets - scalars parameterize complex structure ((4,0) form Ω)
 - $h^{2,1}(X)$ periodic neutral chiral multiplets - periods of C_3 .

M-theory Compactified on CY Fourfold X

- Spectrum ($\mathcal{N} = 2$):
 - $U(1)^{h^{1,1}(X)}$ gauge symmetry - scalars parameterize Kähler class
 $J = \sum x_i e^i$
 - $h^{3,1}(X)$ neutral chiral multiplets - scalars parameterize complex structure ((4,0) form Ω)
 - $h^{2,1}(X)$ periodic neutral chiral multiplets - periods of C_3 .
- Complex and Kähler structure (the latter can be complexified) \leftrightarrow Calabi–Yau metric

M-theory Compactified on CY Fourfold X

- Spectrum ($\mathcal{N} = 2$):
 - $U(1)^{h^{1,1}(X)}$ gauge symmetry - scalars parameterize Kähler class
 $J = \sum x_i e^i$
 - $h^{3,1}(X)$ neutral chiral multiplets - scalars parameterize complex structure ((4,0) form Ω)
 - $h^{2,1}(X)$ periodic neutral chiral multiplets - periods of C_3 .
- Complex and Kähler structure (the latter can be complexified) \leftrightarrow Calabi–Yau metric
- Superpotential **Witten, Sethi Vafa Witten**
 $W = \int_X \Omega \wedge \frac{G}{2\pi}$, $\widetilde{W} = \int_X J \wedge J \wedge \frac{G}{2\pi}$
Critical points: G is (2,2) and primitive.

M-theory Compactified on CY Fourfold X

- Spectrum ($\mathcal{N} = 2$):
 - $U(1)^{h^{1,1}(X)}$ gauge symmetry - scalars parameterize Kähler class
 $J = \sum x_i e^i$
 - $h^{3,1}(X)$ neutral chiral multiplets - scalars parameterize complex structure ((4,0) form Ω)
 - $h^{2,1}(X)$ periodic neutral chiral multiplets - periods of C_3 .
- Complex and Kähler structure (the latter can be complexified) \leftrightarrow Calabi–Yau metric
- Superpotential **Witten, Sethi Vafa Witten**
 $W = \int_X \Omega \wedge \frac{G}{2\pi}$, $\widetilde{W} = \int_X J \wedge J \wedge \frac{G}{2\pi}$
Critical points: G is (2,2) and primitive.

Moduli space of M-theory compactifications is (fibered over) subspace of CY moduli space determined by G

M-theory Compactified on CY Fourfold X

- Spectrum ($\mathcal{N} = 2$):
 - $U(1)^{h^{1,1}(X)}$ gauge symmetry - scalars parameterize Kähler class
 $J = \sum x_i e^i$
 - $h^{3,1}(X)$ neutral chiral multiplets - scalars parameterize complex structure ((4,0) form Ω)
 - $h^{2,1}(X)$ periodic neutral chiral multiplets - periods of C_3 .
- Complex and Kähler structure (the latter can be complexified) \leftrightarrow Calabi–Yau metric
- Superpotential **Witten, Sethi Vafa Witten**
 $W = \int_X \Omega \wedge \frac{G}{2\pi}$, $\widetilde{W} = \int_X J \wedge J \wedge \frac{G}{2\pi}$
Critical points: G is (2,2) and primitive.

Moduli space of M-theory compactifications is (fibered over) subspace of CY moduli space determined by G

- Superpotential can be corrected by **fivebrane instantons Witten**

Constraints on G

- Quantization: $\frac{G}{2\pi} - \frac{c_2(X)}{2} \in H^4(X, \mathbb{Z})$ Witten

Constraints on G

- Quantization: $\frac{G}{2\pi} - \frac{c_2(X)}{2} \in H^4(X, \mathbb{Z})$ **Witten**
- Determined by M2-brane worldvolume anomaly
- Tadpole: **Sethi Vafa Witten**

Constraints on G

- Quantization: $\frac{G}{2\pi} - \frac{c_2(X)}{2} \in H^4(X, \mathbb{Z})$ **Witten**
- Determined by M2-brane worldvolume anomaly
- Tadpole: **Sethi Vafa Witten**

$$M = \frac{\chi(X)}{24} - \frac{1}{2} \int_X \frac{G}{2\pi} \wedge \frac{G}{2\pi} \quad (1)$$

- Gauss' law for G ; $M \geq 0$ is number of space-filling M2-branes.

Constraints on G

- Quantization: $\frac{G}{2\pi} - \frac{c_2(X)}{2} \in H^4(X, \mathbb{Z})$ **Witten**
- Determined by M2-brane worldvolume anomaly
- Tadpole: **Sethi Vafa Witten**

$$M = \frac{\chi(X)}{24} - \frac{1}{2} \int_X \frac{G}{2\pi} \wedge \frac{G}{2\pi} \quad (1)$$

- Gauss' law for G ; $M \geq 0$ is number of space-filling M2-branes.

In general we will not be able to set $G^\sharp = G^\flat = 0$ and $M^\sharp = M^\flat$

The Degeneration

- Consider a codimension- $N - 1$ face σ of Kähler cone of X^\sharp at which a divisor contracts to a **smooth** surface S of transverse A_{N-1} singularities.

The Degeneration

- Consider a codimension- $N - 1$ face σ of Kähler cone of X^\sharp at which a divisor contracts to a **smooth** surface S of transverse A_{N-1} singularities.
- A neighborhood of S in the singular X_0 is

$$xy = z^N \tag{2}$$

in the total space of $\mathcal{O}(\mathcal{L}_1) \oplus \mathcal{O}(\mathcal{L}_2) \oplus K_S$.

- This is resolved to X^\sharp by blowing up $N - 1$ times along S producing $N - 1$ exceptional divisors E_i

The Degeneration

- Consider a codimension- $N - 1$ face σ of Kähler cone of X^\sharp at which a divisor contracts to a **smooth** surface S of transverse A_{N-1} singularities.
- A neighborhood of S in the singular X_0 is

$$xy = z^N \tag{2}$$

in the total space of $\mathcal{O}(\mathcal{L}_1) \oplus \mathcal{O}(\mathcal{L}_2) \oplus K_S$.

- This is resolved to X^\sharp by blowing up $N - 1$ times along S producing $N - 1$ exceptional divisors E_i
- E_i are \mathbb{P}^1 fibrations over S and intersect in common sections according to A_{N-1} Dynkin diagram.

The Degeneration

- Consider a codimension- $N - 1$ face σ of Kähler cone of X^\sharp at which a divisor contracts to a **smooth** surface S of transverse A_{N-1} singularities.
- A neighborhood of S in the singular X_0 is

$$xy = z^N \tag{2}$$

in the total space of $\mathcal{O}(\mathcal{L}_1) \oplus \mathcal{O}(\mathcal{L}_2) \oplus K_S$.

- This is resolved to X^\sharp by blowing up $N - 1$ times along S producing $N - 1$ exceptional divisors E_i
- E_i are \mathbb{P}^1 fibrations over S and intersect in common sections according to A_{N-1} Dynkin diagram.
- Assume as well $G^\sharp = 0$ ensuring that resolution is possible in M-theory

The Degeneration

- Consider a codimension- $N - 1$ face σ of Kähler cone of X^\sharp at which a divisor contracts to a **smooth** surface S of transverse A_{N-1} singularities.
- A neighborhood of S in the singular X_0 is

$$xy = z^N \tag{2}$$

in the total space of $\mathcal{O}(\mathcal{L}_1) \oplus \mathcal{O}(\mathcal{L}_2) \oplus K_S$.

- This is resolved to X^\sharp by blowing up $N - 1$ times along S producing $N - 1$ exceptional divisors E_i
- E_i are \mathbb{P}^1 fibrations over S and intersect in common sections according to A_{N-1} Dynkin diagram.
- Assume as well $G^\sharp = 0$ ensuring that resolution is possible in M-theory
- If $M^\sharp \neq 0$ keep all M2-branes well away from E_i so their worldvolume degrees of freedom decouple

The Degeneration

- Consider a codimension- $N - 1$ face σ of Kähler cone of X^\sharp at which a divisor contracts to a **smooth** surface S of transverse A_{N-1} singularities.
- A neighborhood of S in the singular X_0 is

$$xy = z^N \tag{2}$$

in the total space of $\mathcal{O}(\mathcal{L}_1) \oplus \mathcal{O}(\mathcal{L}_2) \oplus K_S$.

- This is resolved to X^\sharp by blowing up $N - 1$ times along S producing $N - 1$ exceptional divisors E_i
- E_i are \mathbb{P}^1 fibrations over S and intersect in common sections according to A_{N-1} Dynkin diagram.
- Assume as well $G^\sharp = 0$ ensuring that resolution is possible in M-theory
- If $M^\sharp \neq 0$ keep all M2-branes well away from E_i so their worldvolume degrees of freedom decouple

Twisted Reduction on S

- Deep in interior of σ bulk modes decouple and S is large

Twisted Reduction on S

- Deep in interior of σ bulk modes decouple and S is large
- Obtain low-energy theory by twisted reduction $7d \rightarrow 3d$ along S of supersymmetric $SU(N)$ theory.

Twisted Reduction on S

- Deep in interior of σ bulk modes decouple and S is large
- Obtain low-energy theory by twisted reduction $7d \rightarrow 3d$ along S of supersymmetric $SU(N)$ theory.
- Gauge field A_M , triplet S_i of scalars, and a doublet of gaugini Ψ_α satisfying symplectic Majorana condition

Twisted Reduction on S

- Deep in interior of σ bulk modes decouple and S is large
- Obtain low-energy theory by twisted reduction $7d \rightarrow 3d$ along S of supersymmetric $SU(N)$ theory.
- Gauge field A_M , triplet S_i of scalars, and a doublet of gaugini Ψ_α satisfying symplectic Majorana condition
- After twisting and reducing on S SUSY configurations are given by solutions to Vafa–Witten equations

Twisted Reduction on S

- Deep in interior of σ bulk modes decouple and S is large
- Obtain low-energy theory by twisted reduction $7d \rightarrow 3d$ along S of supersymmetric $SU(N)$ theory.
- Gauge field A_M , triplet S_i of scalars, and a doublet of gaugini Ψ_α satisfying symplectic Majorana condition
- After twisting and reducing on S SUSY configurations are given by solutions to Vafa–Witten equations

$$\begin{aligned} F^{(2,0)} = 0, \quad J \wedge F^{(1,1)} + [q, \bar{q}] = 0, \quad [q, \Phi] = 0, \\ D\Phi = \bar{D}\Phi = 0, \quad \bar{D}q = 0; \end{aligned} \quad (3)$$

- F is a connection on an $SU(N)$ bundle over S ; Φ an adjoint-valued real scalar and q an adjoint-valued two-form.

Twisted Reduction on S

- Deep in interior of σ bulk modes decouple and S is large
- Obtain low-energy theory by twisted reduction $7d \rightarrow 3d$ along S of supersymmetric $SU(N)$ theory.
- Gauge field A_M , triplet S_i of scalars, and a doublet of gaugini Ψ_α satisfying symplectic Majorana condition
- After twisting and reducing on S SUSY configurations are given by solutions to Vafa–Witten equations

$$\begin{aligned} F^{(2,0)} = 0, \quad J \wedge F^{(1,1)} + [q, \bar{q}] = 0, \quad [q, \Phi] = 0, \\ D\Phi = \bar{D}\Phi = 0, \quad \bar{D}q = 0; \end{aligned} \quad (3)$$

- F is a connection on an $SU(N)$ bundle over S ; Φ an adjoint-valued real scalar and q an adjoint-valued two-form.

The Low-Energy Gauge Theory

- In our case the bundle is *flat* as a result of $G^\sharp = 0$

The Low-Energy Gauge Theory

- In our case the bundle is *flat* as a result of $G^\sharp = 0$
- Linearize around this to find $\mathcal{N} = 2$ $SU(N)_0$ (A_μ, ϕ) with $h^{2,0}(S) + h^{0,1}(S)$ chiral multiplets q^A, a^i in the adjoint

The Low-Energy Gauge Theory

- In our case the bundle is *flat* as a result of $G^\sharp = 0$
- Linearize around this to find $\mathcal{N} = 2$ $SU(N)_0$ (A_μ, ϕ) with $h^{2,0}(S) + h^{0,1}(S)$ chiral multiplets q^A, a^i in the adjoint
- Cubic superpotential determined by intersection theory on S , but **we restrict to $h^{0,1}(S) = 0$** so this vanishes.

The Low-Energy Gauge Theory

- In our case the bundle is *flat* as a result of $G^\sharp = 0$
- Linearize around this to find $\mathcal{N} = 2$ $SU(N)_0$ (A_μ, ϕ) with $h^{2,0}(S) + h^{0,1}(S)$ chiral multiplets q^A, a^i in the adjoint
- Cubic superpotential determined by intersection theory on S , but **we restrict to $h^{0,1}(S) = 0$** so this vanishes.

The Coulomb Branch

- Coulomb branch: ϕ acquires expectation value, diagonal up to gauge symmetry; eigenvalues permuted by Weyl symmetry.

The Coulomb Branch

- Coulomb branch: ϕ acquires expectation value, diagonal up to gauge symmetry; eigenvalues permuted by Weyl symmetry.
- Gauge symmetry generically $U(1)^{N-1}$. $\rho_g(N-1)$ massless neutral chiral multiplets (complex structure moduli of X^\sharp).

The Coulomb Branch

- Coulomb branch: ϕ acquires expectation value, diagonal up to gauge symmetry; eigenvalues permuted by Weyl symmetry.
- Gauge symmetry generically $U(1)^{N-1}$. $\rho_g(N-1)$ massless neutral chiral multiplets (complex structure moduli of X^\sharp).
- Moduli space is a quotient of the product.

The Coulomb Branch

- Coulomb branch: ϕ acquires expectation value, diagonal up to gauge symmetry; eigenvalues permuted by Weyl symmetry.
- Gauge symmetry generically $U(1)^{N-1}$. $\rho_g(N-1)$ massless neutral chiral multiplets (complex structure moduli of X^\sharp).
- Moduli space is a quotient of the product.

The Higgs Branch

- For $p_g = h^{2,0}(S) > 0$ we also have a Higgs branch

The Higgs Branch

- For $p_g = h^{2,0}(S) > 0$ we also have a Higgs branch
- Gauge symmetry completely broken, ϕ massive.

The Higgs Branch

- For $p_g = h^{2,0}(S) > 0$ we also have a Higgs branch
- Gauge symmetry completely broken, ϕ massive.

$$\dim_{\mathbb{C}} \mathcal{H} = (p_g - 1)(N^2 - 1) . \quad (4)$$

- This is a local model for moduli space of M-theory compactifications on X^b . Meets Coulomb branch along $p_g(N - 1)$ dimensional root.

The Higgs Branch

- For $p_g = h^{2,0}(S) > 0$ we also have a Higgs branch
- Gauge symmetry completely broken, ϕ massive.

$$\dim_{\mathbb{C}} \mathcal{H} = (p_g - 1)(N^2 - 1). \quad (4)$$

- This is a local model for moduli space of M-theory compactifications on X^b . Meets Coulomb branch along $p_g(N - 1)$ dimensional root.
- There are also mixed branches. In codimension $N - p$ in Coulomb branch we have unbroken

$$SU(k_1) \times \cdots \times SU(k_p) \times U(1)^{p-1} \quad (5)$$

- Higgs branch emanating from this has dimension

$$\dim_{\mathbb{C}} \mathcal{H}_{(k_1, \dots, k_p)} = (p_g - 1) \sum_{i=1}^p (k_i^2 - 1) + p_g(p - 1) \quad (6)$$

Transition along Coulomb branch

- We can effect the transition without encountering enhanced gauge symmetry

Transition along Coulomb branch

- We can effect the transition without encountering enhanced gauge symmetry
- Along Coulomb branch turn on generic expectation values for $p_g(N-1)$ neutral chiral multiplets
- Now when $\phi \rightarrow 0$ unbroken gauge symmetry is $U(1)^{N-1}$ and massless charged matter is $(p_g - 1)N(N-1)$ multiplets with charges of adjoints
- When these acquire expectation values Higgsing symmetry completely we find Higgs branch of dimension

$$\dim_{\mathbb{C}} \mathcal{H} = (p_g - 1)N(N-1) - (N-1) + p_g(N-1) = (p_g - 1)(N^2 - 1) \quad (7)$$

The Higgs Branch

- Deforming the singular equation we find the local form for X^b

$$xy = z^N + \sum_{j=0}^N \omega_{N-j} z^j . \quad (8)$$

ω_j are sections of the pluri-canonical line bundles jK_S

The Higgs Branch

- Deforming the singular equation we find the local form for X^b

$$xy = z^N + \sum_{j=0}^N \omega_{N-j} z^j . \quad (8)$$

ω_j are sections of the pluri-canonical line bundles jK_S

- We can compute $\chi(X^b) - \chi(X^\#)$ by using the (singular) form $xy = z^N + \omega_N$

The Higgs Branch

- Deforming the singular equation we find the local form for X^b

$$xy = z^N + \sum_{j=0}^N \omega_{N-j} z^j . \quad (8)$$

ω_j are sections of the pluri-canonical line bundles jK_S

- We can compute $\chi(X^b) - \chi(X^\sharp)$ by using the (singular) form $xy = z^N + \omega_N$
- Given by bouquet of $N - 1$ two-spheres collapsing over vanishing curve $\mathcal{C} \in S$ of ω_N

The Higgs Branch

- Deforming the singular equation we find the local form for X^b

$$xy = z^N + \sum_{j=0}^N \omega_{N-j} z^j . \quad (8)$$

ω_j are sections of the pluri-canonical line bundles jK_S

- We can compute $\chi(X^b) - \chi(X^\sharp)$ by using the (singular) form $xy = z^N + \omega_N$
- Given by bouquet of $N - 1$ two-spheres collapsing over vanishing curve $\mathcal{C} \in S$ of ω_N
- X^\sharp given by resolving along \mathcal{C} so

$$\chi(X^\sharp) - \chi(X^b) = (N - 1)\chi(\mathcal{C}) \quad (9)$$

The Higgs Branch

- Deforming the singular equation we find the local form for X^b

$$xy = z^N + \sum_{j=0}^N \omega_{N-j} z^j . \quad (8)$$

ω_j are sections of the pluri-canonical line bundles jK_S

- We can compute $\chi(X^b) - \chi(X^\sharp)$ by using the (singular) form $xy = z^N + \omega_N$
- Given by bouquet of $N - 1$ two-spheres collapsing over vanishing curve $\mathcal{C} \in S$ of ω_N
- X^\sharp given by resolving along \mathcal{C} so

$$\chi(X^\sharp) - \chi(X^b) = (N - 1)\chi(\mathcal{C}) \quad (9)$$

- Compute $\chi(\mathcal{C}) = N(N + 1) K_S^2$ so

$$\chi(X^b) - \chi(X^\sharp) = N(N - 1)(N + 1)K_S^2 . \quad (10)$$

The Higgs Branch

- Deforming the singular equation we find the local form for X^b

$$xy = z^N + \sum_{j=0}^N \omega_{N-j} z^j . \quad (8)$$

ω_j are sections of the pluri-canonical line bundles jK_S

- We can compute $\chi(X^b) - \chi(X^\sharp)$ by using the (singular) form $xy = z^N + \omega_N$
- Given by bouquet of $N - 1$ two-spheres collapsing over vanishing curve $\mathcal{C} \in S$ of ω_N
- X^\sharp given by resolving along \mathcal{C} so

$$\chi(X^\sharp) - \chi(X^b) = (N - 1)\chi(\mathcal{C}) \quad (9)$$

- Compute $\chi(\mathcal{C}) = N(N + 1) K_S^2$ so

$$\chi(X^b) - \chi(X^\sharp) = N(N - 1)(N + 1)K_S^2 . \quad (10)$$

- In general, Coulomb branch of $\mathcal{N} = 2$ gauge theory suffers quantum corrections

Comments on Quantum Corrections

- In general, Coulomb branch of $\mathcal{N} = 2$ gauge theory suffers quantum corrections
- For $p_g \geq 1$ there will be no corrections to the superpotential

- In general, Coulomb branch of $\mathcal{N} = 2$ gauge theory suffers quantum corrections
- For $p_g \geq 1$ there will be no corrections to the superpotential
- In M-theory fivebrane instantons will be suppressed deep in the cone σ unless they wrap the shrinking divisors E_i . These can generate a superpotential if the divisor has $\chi(E_i, \mathcal{O}_{E_i}) = 1$. Fibration gives $\chi(E_i, \mathcal{O}_{E_i}) = p_g - q + 1$ here so (trivial) agreement. cf. [Diaconescu](#), [Gukov](#)

An Example

- X_0 is CY hypersurface in $\mathbb{P}(1,1,2,2,2,2)$

An Example

- X_0 is CY hypersurface in $\mathbb{P}(1,1,2,2,2,2)$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$

An Example

- X_0 is CY hypersurface in $\mathbb{P}(1,1,2,2,2,2)$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$

An Example

- X_0 is CY hypersurface in $\mathbb{P}(1,1,2,2,2,2)$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$
- With 4 adjoints expect $3 \cdot 3 - 4 = 5$ new deformations

An Example

- X_0 is CY hypersurface in $\mathbb{P}^{(1,1,2,2,2,2)}$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$
- With 4 adjoints expect $3 \cdot 3 - 4 = 5$ new deformations
- To describe X^\flat embed $\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6$ by

$$\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6, \quad (x_1, \dots, x_6) \mapsto (x_1^2, x_2^2, x_1 x_2, x_3, x_4, x_5, x_6) \quad (11)$$

An Example

- X_0 is CY hypersurface in $\mathbb{P}^{(1,1,2,2,2,2)}$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$
- With 4 adjoints expect $3 \cdot 3 - 4 = 5$ new deformations
- To describe X^\flat embed $\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6$ by

$$\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6, \quad (x_1, \dots, x_6) \mapsto (x_1^2, x_2^2, x_1 x_2, x_3, x_4, x_5, x_6) \quad (11)$$

- X_0 is intersection of a quintic and the singular quadric $y_0 y_1 - y_2^2 = 0$

An Example

- X_0 is CY hypersurface in $\mathbb{P}^{(1,1,2,2,2,2)}$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$
- With 4 adjoints expect $3 \cdot 3 - 4 = 5$ new deformations
- To describe X^b embed $\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6$ by

$$\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6, \quad (x_1, \dots, x_6) \mapsto (x_1^2, x_2^2, x_1 x_2, x_3, x_4, x_5, x_6) \quad (11)$$

- X_0 is intersection of a quintic and the singular quadric $y_0 y_1 - y_2^2 = 0$
- Can deform to $X^b = \mathbb{P}^6[2, 5]$ with $h^{1,1} = 1$; $h^{3,1} = 356$; $h^{2,2} = 1472$;
 $\chi = 2190$

An Example

- X_0 is CY hypersurface in $\mathbb{P}^{(1,1,2,2,2,2)}$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$
- With 4 adjoints expect $3 \cdot 3 - 4 = 5$ new deformations
- To describe X^b embed $\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6$ by

$$\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6, \quad (x_1, \dots, x_6) \mapsto (x_1^2, x_2^2, x_1 x_2, x_3, x_4, x_5, x_6) \quad (11)$$

- X_0 is intersection of a quintic and the singular quadric $y_0 y_1 - y_2^2 = 0$
- Can deform to $X^b = \mathbb{P}^6[2, 5]$ with $h^{1,1} = 1$; $h^{3,1} = 356$; $h^{2,2} = 1472$;
 $\chi = 2190$
- Which four-dimensional subspace is consistent in M-theory?

An Example

- X_0 is CY hypersurface in $\mathbb{P}^{(1,1,2,2,2,2)}$
- Singular along quintic surface $S \in \mathbb{P}^3$ with $p_g = 4$
- X^\sharp is toric resolution with $h^{1,1} = 2$; $h^{3,1} = 350$; $h^{2,2} = 1452$;
 $\chi = 2160$
- With 4 adjoints expect $3 \cdot 3 - 4 = 5$ new deformations
- To describe X^b embed $\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6$ by

$$\mathbb{P}^{(1,1,2,2,2,2)} \rightarrow \mathbb{P}^6, \quad (x_1, \dots, x_6) \mapsto (x_1^2, x_2^2, x_1 x_2, x_3, x_4, x_5, x_6) \quad (11)$$

- X_0 is intersection of a quintic and the singular quadric $y_0 y_1 - y_2^2 = 0$
- Can deform to $X^b = \mathbb{P}^6[2, 5]$ with $h^{1,1} = 1$; $h^{3,1} = 356$; $h^{2,2} = 1472$;
 $\chi = 2190$
- Which four-dimensional subspace is consistent in M-theory?
- What is the G^b with which we obtain X^b ?

- For A_{N-1} transitions along smooth surfaces S of general type with $h^1(S) = 0$ and $G^\sharp = 0$ we have found a natural candidate G^\flat satisfying quantization and tadpole conditions which produces Higgs branch of expected dimension and structure

Summary

- For A_{N-1} transitions along smooth surfaces S of general type with $h^1(S) = 0$ and $G^\# = 0$ we have found a natural candidate G^b satisfying quantization and tadpole conditions which produces Higgs branch of expected dimension and structure

Thanks for Listening

Thanks, Dave, for all you've taught us so far. Happy Birthday!!!