

Ph230b, Winter 2011

Problem Set 3 Solution

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Homework 3

Charged pion decay: In order to incorporate charged pions π^\pm in the Fermi theory of weak interactions, it is natural to add a term $f_\pi p_\mu \pi^\pm$ to the charged current J_μ^- (remember that pions couple derivatively). Here, p_μ is the pion momentum and f_π is the pion decay constant. This leads to the following interaction Lagrangian:

$$\mathcal{L}_\pi = \frac{G_F}{\sqrt{2}} f_\pi p_\rho \pi^\pm \bar{\mu} \gamma^\rho (1 - \gamma_5) \nu_\mu + h.c.$$

Show that the total rate for $\pi^- \rightarrow \mu^- \bar{\nu}_\mu$ decay is

$$\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu) = \frac{G_F^2}{8\pi} f_\pi^2 m_{\pi^-} m_\mu^2 \left(1 - \frac{m_\mu^2}{m_{\pi^-}^2}\right)^2$$

where m_{π^-} (resp. m_μ) is the pion (resp. muon) mass.

Kaon decay: Predict the ratio of the $K^- \rightarrow e^- \bar{\nu}_e$ and $K^- \rightarrow \mu^- \bar{\nu}_\mu$ decay rates. Given that the lifetime of the K^- is $\tau = 1.2 \times 10^{-8}$ sec and the $K^- \rightarrow \mu^- \bar{\nu}_\mu$ branching ratio is 64%, estimate the decay constant f_K .

Solution

Charged pion decay

From the Lagrangian we obtain the scattering amplitude;

$$\mathcal{T} = \frac{G_F}{\sqrt{2}} f_\pi p_\rho \bar{u} \gamma^\rho (1 - \gamma_5) v,$$

where \bar{u}, v are four-component spinor of μ^- and $\bar{\nu}_\mu$. From this, we calculate $|\mathcal{T}|^2$, and by summing over final spins, we get

$$\sum_{\text{final spins}} |\mathcal{T}^2| = \frac{G_F^2}{2} f_\pi^2 \text{Tr} \left[(-\not{p}_\mu + m_\mu) \gamma^\rho (1 - \gamma_5) (-\not{p}_{\bar{\nu}_\mu}) \gamma^\delta (1 - \gamma_5) \right] p_\rho p_\delta \quad (1)$$

After some calculation, eq(1) becomes

$$\sum_{\text{final spins}} |\mathcal{T}^2| = 4G_F^2 f_\pi^2 (2(p_\mu \cdot p_{\pi^-})(p_{\pi^-} \cdot p_{\bar{\nu}_\mu}) - (p_\mu \cdot p_{\bar{\nu}_\mu})(p_{\pi^-} \cdot p_{\pi^-})). \quad (2)$$

In the center of mass frame, $p_{\pi^-} = (m_{\pi^-}, 0)$, $p_{\mu} = (E_{\mu}, \mathbf{p})$, and $p_{\bar{\nu}_{\mu}} = (|\mathbf{p}|, -\mathbf{p})$ where $E_{\mu}^2 = m_{\mu}^2 + \mathbf{p}^2$. From $m_{\pi^-} = E_{\mu} + |\mathbf{p}|$, we obtain $|\mathbf{p}| = \frac{m_{\pi^-}^2 - m_{\mu}^2}{2m_{\pi^-}}$, $E_{\mu} = \frac{m_{\pi^-}^2 + m_{\mu}^2}{2m_{\pi^-}}$. Then eq(2) becomes

$$2G_F^2 f_{\pi}^2 m_{\mu}^2 m_{\pi^-}^2 \left(1 - \frac{m_{\mu}^2}{m_{\pi^-}^2}\right).$$

From the differential decay rate,

$$d\Gamma = \frac{1}{2m_{\pi^-}} |\mathcal{T}|^2 \frac{|\mathbf{p}|}{16\pi^2 m_{\pi^-}} d\Omega_{CM},$$

we get

$$\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_{\mu}) = \frac{1}{8\pi} G_F^2 f_{\pi}^2 m_{\mu}^2 m_{\pi^-} \left(1 - \frac{m_{\mu}^2}{m_{\pi^-}^2}\right)^2. \quad (3)$$

Kaon decay

The interaction Lagrangian is similar for $K^- \rightarrow e^- \bar{\nu}_e$ and $K^- \rightarrow \mu^- \bar{\nu}_{\mu}$, so we use eq(3);

$$\begin{aligned} \Gamma(K^- \rightarrow e^- \bar{\nu}_e) &= \frac{1}{8\pi} G_F^2 f_K^2 m_e^2 m_{K^-} \left(1 - \frac{m_e^2}{m_{K^-}^2}\right)^2 \\ \Gamma(K^- \rightarrow \mu^- \bar{\nu}_{\mu}) &= \frac{1}{8\pi} G_F^2 f_K^2 m_{\mu}^2 m_{K^-} \left(1 - \frac{m_{\mu}^2}{m_{K^-}^2}\right)^2 \end{aligned}$$

With $G_F = 1.166 \times 10^{-5} \text{ GeV}^{-2}$, $m_e = 0.511 \text{ MeV}$, $m_{\mu} = 0.106 \text{ GeV}$, $m_{K^-} = 0.494 \text{ GeV}$,

$$\begin{aligned} \Gamma(K^- \rightarrow e^- \bar{\nu}_e) &\simeq 6.978 \times 10^{-28} f_K^2 \text{ eV}^{-1} \\ \Gamma(K^- \rightarrow \mu^- \bar{\nu}_{\mu}) &\simeq 2.732 \times 10^{-23} f_K^2 \text{ eV}^{-1} \end{aligned}$$

Thus, the ratio of the $K^- \rightarrow e^- \bar{\nu}_e$ and $K^- \rightarrow \mu^- \bar{\nu}_{\mu}$ decay rates is

$$\frac{\Gamma(K^- \rightarrow e^- \bar{\nu}_e)}{\Gamma(K^- \rightarrow \mu^- \bar{\nu}_{\mu})} \simeq 2.554 \times 10^{-5},$$

which is close to 2.493×10^{-5} from the PDG.

Given that the lifetime of the K^- is $\tau = 1.2 \times 10^{-8} \text{ sec}$ and the $K^- \rightarrow \mu^- \bar{\nu}_{\mu}$ branching ratio is 64%, we have

$$2.732 \times 10^{-23} f_K^2 \text{ eV}^{-1} \cdot \frac{1}{1.055 \times 10^{-34} \text{ J} \cdot \text{s}} \cdot \frac{1.6 \times 10^{-19} \text{ J}}{\text{eV}} = \frac{1}{1.2 \times 10^{-8} \text{ s}} \frac{64}{100}.$$

Therefore, $f_K \simeq 35.867 \text{ MeV}$. Similar calculation for $\pi^- \rightarrow \mu^- \bar{\nu}_{\mu}$ gives $f_{\pi} \simeq 128.904 \text{ MeV}$.

By the way, from the PDG $f_K \simeq 156.1 \text{ MeV}$, which is a lot different from the above value. This can be explained as follows. Actually, there is a $\sin \theta_C$ in the interaction Lagrangian for $K^- \rightarrow \mu^- \bar{\nu}_{\mu}$ and $\cos \theta_C$ for $\pi^- \rightarrow \mu^- \bar{\nu}_{\mu}$ where θ_C is the Cabibbo angle. These factors come from the coupling of \bar{u} quark to $d' = d \cos \theta_C + s \sin \theta_C$ when considering the full theory. Generalization of this to all quarks is encoded in the CKM matrix. Considering these, with $f_{K,\text{PDG}} = 156.1 \text{ MeV}$ and $f_{\pi,\text{PDG}} = 130.41 \text{ MeV}$, we get $f_{K,\text{here}} \simeq \sin \theta_C f_{K,\text{PDG}} = 35.358 \text{ MeV}$ and $f_{\pi,\text{here}} \simeq \cos \theta_C f_{\pi,\text{PDG}} = 127.016 \text{ MeV}$. Above calculated results, $f_{K,\text{calculated here}} \simeq 35.867 \text{ MeV}$ and $f_{\pi,\text{calculated here}} \simeq 128.904 \text{ MeV}$ are close to these values.