

### Week 3 (due Oct. 23)

Reading: Srednicki, sections 4 and 5.

1. Compute the commutator function  $[\phi(x), \phi(0)]$  for the free real scalar field  $\phi$  with zero mass ( $m = 0$ ). Hint: by rotational invariance, you may assume that the spatial part of  $x$  is along the  $x^1$  axis. The integral over  $k_2$  and  $k_3$  is easily computed, if we recall that

$$\frac{d^3k}{2\omega_k} = d^4k \delta(-k^2)\theta(k^0).$$

. The remaining integral over  $k^0$  and  $k^1$  is most easily evaluated in the “light-cone coordinates”  $k_+ = k^0 - k^1$  and  $k_- = k^0 + k^1$ .

2. Let  $\phi$  be as in problem 1. Compute the vacuum expectation value

$$\langle 0|\phi(x)\phi(0)|0 \rangle .$$

Hint: be careful, this is a distribution, not a function. Use the same method as in problem 1.

3. The Hamiltonian for the free complex scalar field of mass  $m$  is

$$H = \int d^3x (p^\dagger p + \partial_i \phi^\dagger \partial_i \phi + m^2 \phi^\dagger \phi) .$$

Here  $p = \partial_0 \phi^\dagger$  is the momentum conjugate to  $\phi$  and  $p^\dagger = \partial_0 \phi$  is the momentum conjugate to  $\phi^\dagger$ . The nonvanishing equal-time commutators are

$$[p(\vec{x}), \phi(\vec{y})] = -i\delta^3(\vec{x} - \vec{y}), \quad [p^\dagger(\vec{x}), \phi^\dagger(\vec{y})] = -i\delta^3(\vec{x} - \vec{y}).$$

Show that the Heisenberg equations of motion

$$i\partial_0 \phi = [H, \phi], \quad i\partial_0 p = [H, p]$$

are equivalent to the Klein-Gordon equation for  $\phi$ .

4. (a) Consider a field theory with three real scalar fields  $\phi^a(x)$ ,  $a = 1, 2, 3$ , and a Lagrangian

$$\mathcal{L} = -\frac{1}{2} \partial_\mu \phi^a(x) \partial^\mu \phi^a(x) - V(\phi^a \phi^a).$$

Here summation over repeating indices  $a$  is assumed, and  $V$  is an arbitrary function. This Lagrangian is obviously invariant with respect to orthogonal transformations of the fields  $\phi^a$ :

$$\phi^a(x) \mapsto \tilde{\phi}^a(x) = R_b^a \phi^b(x),$$

where  $R_b^a$  is a constant orthogonal  $3 \times 3$  matrix. The rotation group in three dimensional space has dimension three, so we expect to get three conserved currents. Show that infinitesimal transformations for  $\phi^a(x)$  can be put into the form

$$\delta\phi^a(x) = \epsilon^{abc}\phi^b(x)\beta^c,$$

where  $\beta^c, c = 1, 2, 3$  parametrize an infinitesimal rotation, and  $\epsilon^{abc}$  is a completely anti-symmetric tensor uniquely defined by the condition  $\epsilon^{123} = 1$ . Deduce the conserved currents corresponding to this symmetry.

(b) Let the currents found in part (a) be called  $J^{a\mu}, a = 1, 2, 3$ . The corresponding charges are

$$Q^a = \int d^3x J^{a0}(x).$$

Compute the commutator of  $Q^a$  and  $Q^b$  using canonical commutation relations for  $\phi^a$  and their time derivatives. Show that  $Q^a$  form a Lie algebra isomorphic to the Lie algebra of the rotation group (i.e. show that they obey the same commutation relations as components of the angular momentum operator in quantum mechanics).