Reading: Srednicki sec. 22, pp. 132-135.

1. Consider the theory of a complex scalar field  $\phi$ . Let  $j^{\mu} = i(\phi^{\dagger}\partial^{\mu}\phi - (\partial^{\mu}\phi^{\dagger})\phi)$ . This is a conserved current, and the corresponding conserved charge is

$$Q = \int j^0 d^3x.$$

Express Q in terms of creation and annihilation operators  $a, a^{\dagger}, b, b^{\dagger}$  and show that Q is the number of particles minus the number of antiparticles.

2. The real scalar field has action

$$S = \int d^4x \left( -\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2 \right).$$

It is natural to define the operator of spatial momentum for the free real scalar field by a Fock-space expression

$$\mathbf{P} = \int \frac{d^3k}{(2\pi)^3 2\omega_{\mathbf{k}}} \,\mathbf{k} \, a_k^{\dagger} a_k.$$

(a) Show that **P** satisfies

$$[P^{j}, \phi(x)] = i\partial_{j}\phi(x), \quad j = 1, 2, 3.$$

- (b) Express **P** in terms of  $\phi(x)$  and  $p(x) = \partial_0 \phi(x)$ .
- 3. (a) Consider the real scalar field with m=0. Show that the following transformation is a symmetry of the action:

$$\tilde{\phi}(x) = \lambda^{-1}\phi(\lambda^{-1}x), \quad \lambda \neq 0.$$

This transformation is called a dilatation. An alternative way to write this transformation is to let  $\tilde{x} = \lambda x$  and write

$$\tilde{\phi}(\tilde{x}) = \lambda^{-1}\phi(x).$$

From this viewpoint, it is clear that dilatation symmetry originates from a geometric transformation  $x \to \tilde{x} = \lambda x$ .

Derive the conserved current corresponding to this symmetry and check by an explicit computation that it is conserved if the equations of motion are satisfied. (b) Consider a massless scalar in n-dimensional space-time with the action

$$S_0 = -\frac{1}{2} \int d^n x \, \partial_\mu \phi \partial^\mu \phi.$$

Consider a transformation  $x \to \tilde{x} = \lambda x$ , where  $\lambda$  is a nonzero real number. Find the transformation law for  $\phi(x)$  which makes the action invariant and compute the corresponding conserved current. Further, add the following term to the action:

$$S_1 = -\int d^n x \ V(\phi(x)),$$

where  $V(\phi)$  is a function. What is the most general form of the function  $V(\phi)$  compatible with the dilatation invariance? Compute the dilatation current for such  $V(\phi)$ .

4. (a) Consider a field theory with three real scalar fields  $\phi^a(x)$ , a=1,2,3, and a Lagrangian

$$\mathcal{L} = -\frac{1}{2}\partial_{\mu}\phi^{a}(x)\partial^{\mu}\phi^{a}(x) - V(\phi^{a}\phi^{a}).$$

Here summation over repeating indices a is assumed, and V is an arbitrary function. This Lagrangian is obviously invariant with respect to orthogonal transformations of the fields  $\phi^a$ :

$$\phi^a(x) \mapsto \tilde{\phi}^a(x) = R_b^a \phi^b(x),$$

where  $R_b^a$  is a constant orthogonal  $3 \times 3$  matrix. The rotation group in three dimensional space has dimension three, so we expect to get three conserved currents. Show that infintesimal transformations for  $\phi^a(x)$  can be put into the form

$$\delta \phi^a(x) = \epsilon^{abc} \phi^b(x) \beta^c,$$

where  $\beta^c$ , c=1,2,3 parametrize an infintesimal rotation, and  $\epsilon^{abc}$  is a completely anti-symmetric tensor uniquely defined by the condition  $\epsilon^{123}=1$ . Deduce the conserved currents corresponding to this symmetry.

(b) Let the currents found in part (a) be called  $J^{a\mu}$ , a=1,2,3. The corresponding charges are

$$Q^a = \int d^3x J^{a0}(x).$$

Compute the commutator of  $Q^a$  and  $Q^b$  using canonical commutation relations for  $\phi^a$  and their time derivatives. Show that  $Q^a$  form a Lie algebra isomorphic to the Lie algebra of the rotation group (i.e. show that they obey the same commutation relations as components of the angular momentum operator in quantum mechanics).