

## Ph 12b

### Homework Assignment No. 8 Due: 5pm, Thursday, 11 March 2010

**1. A barrier in a well** (10 points).

A free quantum-mechanical particle with mass  $m$  moves inside a one-dimensional box with impenetrable walls located at  $x = \pm a$ . Furthermore, a *repulsive*  $\delta$ -function barrier sits at the center of the well, so the potential energy function  $V(x)$  in between the two impenetrable walls is given by

$$V(x) = \frac{\hbar^2}{m} \Delta \delta(x) ,$$

where  $\delta(x)$  denotes the Dirac  $\delta$ -function and  $\Delta > 0$ . As explained in Problem 2 last week, this  $\delta$ -function potential causes the logarithmic derivative of the wave function  $\varphi(x)$  to jump discontinuously at the origin:

$$\varphi'(0^+) - \varphi'(0^-) = 2\Delta\varphi(0) .$$

Here  $\varphi'(x)$  denotes the first derivative of  $\varphi(x)$ , and  $\varphi'(0^+)$  (respectively  $\varphi'(0^-)$ ) denotes the limit of  $\varphi'(x)$  as  $x$  approaches zero from positive (negative) values. The sign convention used here for  $\Delta$  is the opposite of that used last week; now  $\Delta > 0$  is the case of a repulsive barrier.

- a) For the even energy eigenstates, what is the value of the logarithmic derivative  $\varphi'(x)/\varphi(x)$  at  $x = 0^+$  and  $x = 0^-$ ?
- b) For the even energy eigenstates, derive an equation that determines the wavenumber  $k$  implicitly, where  $E = \hbar^2 k^2 / 2m$ . Express your answer in the form

$$\Delta a = f(ka),$$

where  $f$  is a suitable function.

- c) Consider the limiting case of an infinitely strong repulsive barrier:  $\Delta a \rightarrow \infty$ . What are values of the energy eigenvalues in this limit, for both even and odd  $n$ ?
- d) Draw rough sketches of the wave functions for the ground state and the first excited state in the limit  $\Delta a \rightarrow \infty$ .

**2. Reflectionless potential** (15 points).

Consider a particle with mass  $m$  moving in the attractive potential

$$V(x) = -\frac{\hbar^2 k_0^2}{m} \operatorname{sech}^2(k_0 x),$$

where  $\operatorname{sech}(z) = 2(e^z + e^{-z})^{-1}$  denotes the hyperbolic secant function.

- a) Show that the time-independent Schrödinger equation for this potential can be expressed as

$$\left( -\frac{d^2}{dz^2} - 2 \operatorname{sech}^2(z) \right) \varphi(z) = \bar{k}^2 \varphi(z), \quad (1)$$

where  $z = k_0 x$  is a dimensionless position variable, and  $\bar{k}^2 k^2 / k_0^2 = 2mE / \hbar^2 k_0^2$  is a dimensionless wavenumber.

- b) Show that

$$(i\bar{k} - \tanh(z))e^{i\bar{k}z} \quad (2)$$

solves eq.(1).

- c) Show that eq.(2) approaches  $Ae^{i\bar{k}z}$  as  $z \rightarrow -\infty$  and approaches  $Ce^{i\bar{k}z}$  as  $z \rightarrow +\infty$ , where  $A$  and  $C$  are constants. What are the values of these constants.
- d) What is the transmission amplitude  $C/A$ ? Show that the transmission probability  $T = |C/A|^2$  is one, and that the reflection probability  $R = 1 - T$  is zero. Hence, if a wavepacket is incident on this potential from the far left, there is no reflected wave packet at all.
- e) Find an imaginary value  $\bar{k} = i\bar{\kappa}$  such that  $A/C = 0$  and the transmission amplitude thus diverges. For this value of  $\bar{\kappa}$ , there is a normalizable bound state solution, which decays exponential for both  $z \rightarrow +\infty$  and  $z \rightarrow -\infty$ .
- f) Check that

$$\varphi(z) = \operatorname{sech}(z)$$

solves eq.(1), where  $\bar{k}^2 = -\bar{\kappa}^2$  and  $\bar{\kappa}$  is the value found in (e). This is the bound state solution. What is the corresponding bound state energy?

**3. Bound states in a linear potential** (15 points).

Consider a particle with mass  $m$  moving in the potential

$$V(x) = F|x|$$

where  $|x|$  denotes the absolute value function. Thus there is a constant force  $F$  directed toward the origin.

a) Show that the time-independent Schrödinger equation for this potential can be expressed in the form (for  $x \geq 0$ )

$$\left(-\frac{d^2}{dy^2} + y\right) \varphi(y) = \bar{E} \varphi(y), \quad (3)$$

where

$$y = \left(\frac{\hbar^2}{2mF}\right)^{-1/3} x, \quad \bar{E} = \left(\frac{\hbar^2 F^2}{2m}\right)^{-1/3} E.$$

Equivalently, we may write

$$\frac{d^2}{dz^2} \varphi(z) = z \varphi(z)$$

where  $z = y - \bar{E}$ . The solution to this equation that decays as  $z \rightarrow +\infty$  is the Airy function  $\text{Ai}(z)$ .

All real zeros of  $\text{Ai}(z)$  and of its first derivative  $\text{Ai}'(z)$  occur for  $z < 0$ . We denote the zeros of  $\text{Ai}'(z)$ , in order of increasing absolute value, by  $a_0, a_2, a_4, \dots$ , and we denote the zeros of  $\text{Ai}(z)$  in order of increasing absolute value by  $a_1, a_3, a_5, \dots$ . These constants have the numerical values:

$$\begin{aligned} -a_0 &= 1.0188 \dots \\ -a_1 &= 2.3381 \dots \\ -a_2 &= 3.2482 \dots \\ -a_3 &= 4.0879 \dots \\ -a_4 &= 4.8201 \dots \\ -a_5 &= 5.5206 \dots \end{aligned} \quad (4)$$

- b) Show that for  $n = 0, 1, 2, \dots$  there is a bound state solution to the Schrödinger equation with  $n$  nodes and dimensionless “energy”  $\bar{E} = \bar{E}_n = -a_n$ .

Using the WKB approximation and the connection formulas, we can derive the Bohr-Sommerfeld criterion:

$$\int_{x_1}^{x_2} dx k(x) = \pi \left( n + \frac{1}{2} \right),$$

where  $n$  is the number of nodes in the bound state wavefunction,  $E_n$  is the corresponding energy,  $x_1$  and  $x_2$  are the classical turning points for  $E = E_n$ , and  $k(x)^2 = 2m(E_n - V(x))/\hbar^2$ . For the harmonic potential, this WKB estimate actually agrees with the exact value of  $E_n$ , but in general there are corrections higher order in  $1/n$ .

- c) Apply the WKB criterion to the linear potential, deriving a formula for  $\bar{E}_n$ . For  $n = 0, 1, 2, 3, 4, 5$ , compare to the exact result from (b). You should find pretty good agreement for all  $n \geq 1$ . Furthermore, you should find (considering the odd and even values of  $n$  separately), that the agreement gets systematically better as  $n$  increases.