

Homework 8 Solutions

Ph 12b Winter 2010

March 13, 2010

1. A Barrier in a Well

- a. For an even energy eigenstate ϕ , $\phi(x) = \phi(-x)$ and $\phi'(x) = -\phi'(-x)$. functions, From the previous problem set, we know

$$\lim_{\epsilon \rightarrow 0} \phi(x - \epsilon) - \phi(x + \epsilon) \Rightarrow \phi(0^-) - \phi(0^+) = 0 \Rightarrow \phi(0^-) = \phi(0^+) \Rightarrow \phi(0^-), \phi(0^+) \rightarrow \phi(0)$$

The derivative matching condition gives us

$$\phi'(0^+) - \phi'(0^-) = 2\Delta\phi(0) \Rightarrow 2\phi'(0^+) = 2\Delta\phi(0) \Rightarrow \frac{\phi'(0^+)}{\phi(0^+)} = \Delta$$

Similarly,

$$\frac{\phi'(0^-)}{\phi(0^-)} = -\Delta.$$

- b. Assume the same functional form of ϕ as last time and use the boundary conditions $\phi(a) = \phi(-a) = 0$ to get

$$\begin{aligned} \phi(x) = \phi(-x) &\Rightarrow Ce^{ikx} + De^{-ikx} = Ae^{-ikx} + Be^{ikx} \Rightarrow A = D, B = C \\ \Rightarrow \phi'(0^+) - \phi'(0^-) &= ikC - ikD - ikA + ikB = 2ik(B - A) = 2\Delta(A + B) = 2\Delta\phi(0) \end{aligned}$$

$$\Rightarrow B = \frac{ik + \Delta}{ik - \Delta} A$$

$$\phi(-a) = Ae^{-ika} + Be^{ika} = Ae^{-ika} + \frac{ik + \Delta}{ik - \Delta} Ae^{ika} = 0$$

$$\begin{aligned} ik(e^{-ika} + e^{ika}) - \Delta(e^{-ika} - e^{ika}) &= 0 \Rightarrow 2ik \cos(ka) + 2i\Delta \sin(ka) \Rightarrow \Delta = -k \cot(ka) \\ \Rightarrow \Delta a &= -ka \cot(ka) \end{aligned}$$

- c. For n odd, the delta barrier does not affect the wavefunction since $\phi(0) = 0$. As $\Delta a \rightarrow \infty$, $\sin(ka) \rightarrow 0$, meaning that $k = o\pi/a$. Let o be the number of nodes; then $o = n + 1/2$ and

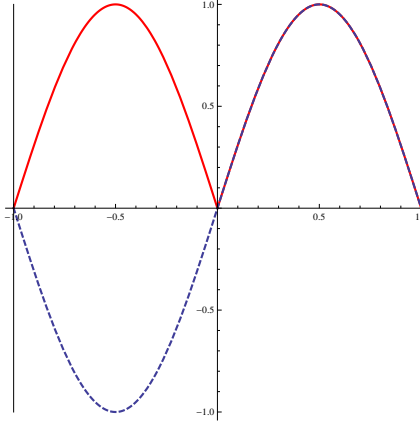
$$E_n = \frac{\hbar^2 \pi^2}{2ma^2} \left(\frac{n+1}{2} \right)^2.$$

For n even, the barrier has an effect at $x = 0$ which requires the creation of an additional node, requiring that $o = n + 2/2$ and

$$E_n = \frac{\hbar^2 \pi^2}{2ma^2} \left(\frac{n+2}{2} \right)^2$$

resulting in degeneracy.

- d. The ground state must be even and is shown in red and the first excited state is in blue.



2. Reflectionless Potential

a.

$$\begin{aligned} \hat{H}\psi(x) = E\psi(x) &\Rightarrow \left[\frac{\hat{p}}{2m} + V(x) \right] \psi(k_0x) = E\psi(k_0x) \\ &\Rightarrow \left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - \frac{\hbar^2}{m} k_0^2 \operatorname{sech}^2(k_0x) \right] \phi(k_0x) = E\phi(k_0x) \quad z = k_0x \Rightarrow dz = k_0 dx \\ &\Rightarrow \left[-\frac{\hbar^2}{2m} \frac{d^2}{dz^2} k_0^2 - \frac{\hbar^2}{m} k_0^2 \operatorname{sech}^2(z) \right] \phi(z) = E\phi(z) \Rightarrow \left[-\frac{d^2}{dz^2} - 2 \operatorname{sech}^2(z) \right] \phi(z) = \underbrace{\frac{2mE}{\hbar^2 k_0^2}}_{\bar{k}^2} \phi(z) \end{aligned}$$

b.

$$\begin{aligned} \left[-\frac{d^2}{dz^2} - 2 \operatorname{sech}^2(z) \right] (i\bar{k} - \tanh z) e^{i\bar{k}z} &= -\frac{d}{dz} \left[-\operatorname{sech}^2 z e^{i\bar{k}z} + i\bar{k} \phi(z) e^{i\bar{k}z} \right] - 2 \operatorname{sech}^2 z \phi(z) \\ &\Rightarrow (\bar{k}^2 + 2 \operatorname{sech}^2 z) \phi(z) - 2 \operatorname{sech}^2 z \phi(z) = \bar{k}^2 \phi(z) \end{aligned}$$

c. We only have to show the limiting behavior of $i\bar{k} - \tanh z$:

$$\lim_{z \rightarrow \infty} i\bar{k} - \tanh z = i\bar{k} - 1 = C \quad \lim_{z \rightarrow -\infty} i\bar{k} - \tanh z = i\bar{k} + 1 = A$$

d.

$$\begin{aligned} \frac{C}{A} &= \frac{i\bar{k} - 1}{i\bar{k} + 1} = \frac{(i\bar{k} - 1)(i\bar{k} - 1)}{(i\bar{k} + 1)(i\bar{k} - 1)} = \frac{\bar{k}^2 + 2i\bar{k} - 1}{\bar{k}^2 + 1} \\ T &= \left| \frac{C}{A} \right|^2 = \frac{1}{(\bar{k}^2 + 1)^2} (\bar{k}^2 + 2i\bar{k} - 1)(\bar{k}^2 - 2i\bar{k} - 1) = \frac{\bar{k}^4 + 2\bar{k}^2 + 1}{(\bar{k}^2 + 1)^2} = 1 \\ R &= 1 - T = 0 \end{aligned}$$

e. Solve:

$$\frac{A}{C} = \frac{i\bar{k} + 1}{i\bar{k} - 1} = \frac{i(i\bar{k}) + 1}{i(i\bar{k}) - 1} = \frac{-\bar{k} + 1}{-\bar{k} - 1} = 0 \Rightarrow \bar{k} = 1 \Rightarrow \bar{k} = i$$

f. $\bar{k}^2 = -\bar{\kappa}^2 = -1$

$$\begin{aligned} \left[-\frac{d^2}{dz^2} - 2 \operatorname{sech}^2(z) \right] \phi(z) &= \frac{d}{dz} \operatorname{sech} z \tanh z - 2 \operatorname{sech}^3 z = \operatorname{sech}^3 z - \operatorname{sech} z \tanh^2 z - 2 \operatorname{sech}^3 z \\ &\Rightarrow -\operatorname{sech} z [\operatorname{sech}^2 z + \tanh^2 z] = -\operatorname{sech} z = \bar{k}^2 \phi(z) \\ \bar{k}^2 &= \frac{2mE}{\hbar^2 k_0^2} \Rightarrow E = -\frac{\hbar^2 k_0^2}{2m} \end{aligned}$$

3. Bound States in a Linear Potential

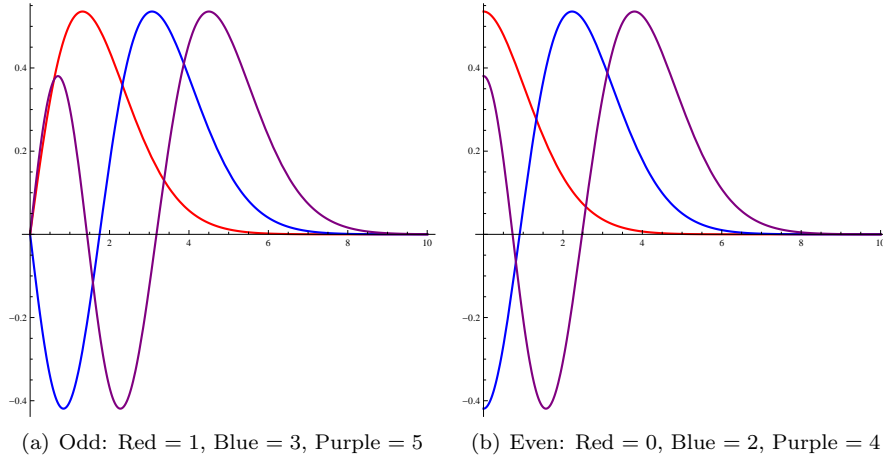
a.

$$\begin{aligned} \hat{H}\psi(x) = E\psi(x) &\Rightarrow \left[-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} - F|x| \right] \phi(x) = E\phi(x) \quad \text{for } x \geq 0 \text{ and } y = \left(\frac{\hbar^2}{2mF} \right)^{-1/3} x \\ \Rightarrow dy &= \left(\frac{\hbar^2}{2mF} \right)^{-1/3} dx \Rightarrow \left[-\frac{\hbar^2}{2m} \frac{d^2}{dy^2} \left(\frac{2mF}{\hbar^2} \right)^{2/3} - F \left(\frac{\hbar^2}{2mF} \right)^{1/3} y \right] \phi(y) = E\phi(y) \\ &\Rightarrow \left(\frac{\hbar^2 F^2}{2m} \right)^{1/3} \left[-\frac{d^2}{dy^2} + y \right] \phi(y) = E\phi(y) \Rightarrow \left[-\frac{d^2}{dy^2} + y \right] \phi(y) = \bar{E}\phi(y) \end{aligned}$$

b. For an even solution $\phi'(0) = 0$ and for an odd solution $\phi(0) = 0$. Then we get

$$\phi(x) \sim \operatorname{Ai}(\alpha x + \bar{E}) \Rightarrow \phi(0) = \begin{cases} \operatorname{Ai}'(\bar{E}) = 0 & \text{even solutions} \\ \operatorname{Ai}(\bar{E}) = 0 & \text{odd solutions} \end{cases}$$

We can see how the Airy solutions look like when shifted by the zeros: When n is odd, there are



$o = n - 1/2$ zeros on the right side excluding the one at the origin (see above plots). Then by symmetry, the whole wavefunction has $2o + 1 = n$ nodes. When n is even, there are $o = n/2$ zeros on the right side. By symmetry, the whole wavefunction has $2o = n$ nodes; thus the solutions have n nodes.

c.

$$\int_{x_1}^{x_2} k(x) dx = \pi \left(n + \frac{1}{2} \right) \quad \text{classical turning points: } x_i = \pm \frac{E_n}{F}$$

$$\begin{aligned}
k(x)^2 &= \frac{2m}{\hbar^2}(E_n - V(x)) \Rightarrow k(x) = \sqrt{\frac{2m}{\hbar^2}(E_n - Fx)} \\
\Rightarrow 2\sqrt{\frac{2m}{\hbar^2}} \int_0^{E_n/F} \sqrt{E_n - Fx} dx &= 2\sqrt{\frac{2m}{\hbar^2}} \left(-\frac{2}{3F}(E_n - Fx)^{3/2} \right) \Big|_0^{E_n/F} = \frac{4}{3F} \sqrt{\frac{2m}{\hbar^2}} E_n^{3/2} = \pi \left(n + \frac{1}{2} \right) \\
\Rightarrow E_n &= \left(\frac{\hbar^2 F^2}{2m} \frac{9}{16} \pi^2 \left(n + \frac{1}{2} \right)^2 \right)^{1/3} = \bar{E}_n = \left(\frac{3}{4} \pi \left(n + \frac{1}{2} \right) \right)^{2/3}
\end{aligned}$$

We then get

n	WKB Prediction	Airy Exact	% difference
0	1.115	1.0188	+9.4%
1	2.320	2.3381	-0.8%
2	3.261	3.2482	+0.4%
3	4.082	4.0879	-0.14%
4	4.827	4.8201	+0.14%
5	5.517	5.5206	-0.06%

At $n = 19$, the 10th of $\text{Ai}(z)$ is 12.8288. The predicted value of $\bar{E}_n = 12.8281$, which is lower by 5×10^{-5} .