

# Ph 219/CS 219

## Exercises

Due: Wednesday 4 March 2009

### 6.1 Finding a collision

Suppose that a black box evaluates a function

$$f : \{0, 1\}^n \rightarrow \{0, 1\}^{n-1} . \quad (1)$$

We are promised that the function is 2-to-1, and we are to find a “collision” – values  $x$  and  $y$  such that  $f(x) = f(y)$ . This problem is harder than Simon’s problem, because we are not promised that the function is periodic. Let  $N = 2^n$ .

- Describe a randomized classical algorithm that requires  $\text{SPACE} = O(\sqrt{N})$  and that succeeds in finding a collision with high probability in  $O(\sqrt{N})$  queries of the black box.
- Now suppose that only  $\text{SPACE} = O(N^{1/3})$  is available. Describe a randomized classical algorithm that finds a collision with high probability in  $O(N^{2/3})$  queries.
- Show that Grover’s exhaustive search algorithm can be used to find a collision in  $O(\sqrt{N})$  quantum queries, using  $\text{SPACE} = O(1)$ .
- Describe a quantum algorithm that uses  $\text{SPACE} = O(M)$  and finds a collision in  $O(M) + O(\sqrt{N/M})$  quantum queries. **[Hint:** First query the box  $M$  times to learn the value of  $f(x)$  for  $M$  arguments  $\{x_1, x_2, \dots, x_M\}$ , then search for  $y$  such that  $f(y) = f(x_i)$  for some  $x_i$ .] Thus, if  $M$  is chosen to optimize the number of queries, the quantum algorithm uses  $\text{SPACE} = O(N^{1/3})$  and  $O(N^{1/3})$  quantum queries.

### 6.2 All the information for half the price

A black box computes a function

$$f : \{0, 1\}^n \rightarrow \{0, 1\} . \quad (2)$$

This function can be represented by a binary string

$$X = X_{N-1}X_{N-2} \cdots X_1X_0 , \quad (3)$$

where  $X_i = f(i)$  and  $N = 2^n$ . Our goal is to obtain, with high probability of success, *complete* information about the box; that is, to find the value of  $X$ . The only resource we care about is the number of queries of the box — TIME and SPACE are otherwise unlimited.

- a) How many classical queries are needed to find  $X$  with success probability at least  $2/3$ ?
- b) Suppose that the state

$$|\Psi_{X,N}\rangle = \frac{1}{\sqrt{2^N}} \sum_{Y \in \{0,1\}^N} (-1)^{X \cdot Y} |Y\rangle \quad (4)$$

has been prepared, where the sum is over all  $N$ -bit strings, and  $X \cdot Y$  denotes the mod 2 bitwise inner product

$$\begin{aligned} X \cdot Y &= (X_{N-1} \cdot Y_{N-1}) \oplus (X_{N-2} \cdot Y_{N-2}) \\ &\quad \cdots \oplus (X_1 \cdot Y_1) \oplus (X_0 \cdot Y_0). \end{aligned} \quad (5)$$

Describe a way, by applying a simple unitary and then a measurement, to find the value of  $X$  with certainty.

- c) Explain how the unitary transformation

$$U : |Y\rangle \rightarrow (-1)^{X \cdot Y} |Y\rangle \quad (6)$$

can be implemented with  $|Y|$  queries of the box, where  $|Y|$  denotes the *Hamming weight* of  $Y$ , the number of 1's in the string.

- d) Suppose we prepare the state

$$|\Phi_K\rangle = \frac{1}{\sqrt{M_K}} \sum_{Y:|Y| \leq K} |Y\rangle, \quad (7)$$

where

$$M_K = \sum_{j=0}^K \binom{N}{j}, \quad (8)$$

and then apply  $U$  (requiring at most  $K$  queries) to obtain

$$|\Psi_{X,K}\rangle = \frac{1}{\sqrt{M_K}} \sum_{Y:|Y| \leq K} (-1)^{X \cdot Y} |Y\rangle, \quad (9)$$

Show that, by applying the procedure that you described in your answer to (b), we can determine the value of  $X$  with a probability of success

$$P_{\text{succ}}(N, K) = |\langle \Psi_{X,K} | \Psi_{X,N} \rangle|^2, \quad (10)$$

and compute the value of  $P_{\text{succ}}(N, K)$ .

e) Suppose that

$$K = N/2 + c\sqrt{N}, \quad (11)$$

where  $c$  is a constant. Show that

$$1 - P_{\text{succ}}(N, K) = O(e^{-2c^2}). \quad (12)$$

Thus we can extract all the information from the box in a number of queries  $(N/2) \cdot [1 + O(1/\sqrt{N})]$ .

### 6.3 Quantum counting

A black box computes a function

$$f : \{0, 1\}^n \rightarrow \{0, 1\}, \quad (13)$$

which can be represented by a binary string

$$X = X_{N-1}X_{N-2} \cdots X_1X_0, \quad (14)$$

where  $X_i = f(i)$  and  $N = 2^n$ . Our goal is to count the number  $r$  of states “marked” by the box; that is, to determine the Hamming weight  $r = |X|$  of  $X$ . We can devise a quantum algorithm that counts the marked states by combining Grover’s exhaustive search with the quantum Fourier transform.

a) Suppose we can consult a quantum oracle that executes the unitary transformation  $U$ . We’d like to perform  $\Lambda(U)$ , the unitary  $U$  conditioned on the value of a control qubit. Devise a quantum circuit with one oracle query that executes  $\Lambda(U)$ , using ancilla qubits and  $\Lambda(\text{SWAP})$  gates, where

$$\text{SWAP} : |x\rangle|y\rangle \rightarrow |y\rangle|x\rangle. \quad (15)$$

b) Let

$$|\Psi_X\rangle = \frac{1}{\sqrt{r}} \sum_{j: X_j=1} |j\rangle \quad (16)$$

denote the uniform superposition of the marked states, and let  $U_{\text{Grover}}$  denote the “Grover iteration,” which performs a rotation by the angle  $2\theta$  in the plane spanned by  $|\Psi_X\rangle$  and

$$|s\rangle = \frac{1}{\sqrt{N}} \sum_{j=0}^N |j\rangle, \quad (17)$$

where

$$\sin \theta = \langle s | \Psi_X \rangle = \sqrt{\frac{r}{N}}. \quad (18)$$

Consider a unitary transformation

$$V : |t\rangle \otimes |\Phi\rangle \rightarrow |t\rangle \otimes U_{\text{Grover}}^t |\Phi\rangle \quad (19)$$

that reads a counter register taking values  $t \in \{0, 1, 2, \dots, T-1\}$  (where  $T = 2^m$ ), and then applies  $U_{\text{Grover}}$   $t$  times. Explain how  $V$  can be implemented, calling the oracle  $T-1$  times. [**Hint:** Use the binary expansion  $t = \sum_{k=0}^{m-1} t_k 2^k$  and the conditional oracle call from (a).]

- c) Suppose that  $r \ll N$ . Show that, by applying  $V$ , performing the quantum Fourier transform on the counter register, and then measuring the counter register, we can determine  $\theta$  to accuracy  $O(1/T)$ , and hence we can find  $r$  with high success probability in  $T = O(\sqrt{rN})$  queries. Compare to the best classical protocol.

#### 6.4 Simulating the Schrödinger equation

A quantum computer can simulate the continuous time evolution of a quantum system governed by a “geometrically local” Hamiltonian. Usually such simulations are done by approximating continuous time evolution with a series of discrete time steps, introducing an error that should be small if the step size is small enough.

For example, by expanding the exponential in a power series, you can verify that

$$e^{A+B} - e^A e^B = -\frac{1}{2}[A, B] + \dots; \quad (20)$$

here  $[A, B] = AB - BA$  is the commutator of operators  $A$  and  $B$  and the ellipsis indicates terms that are higher order in  $A$  and  $B$ . Suppose the (time-independent) Hamiltonian  $H$  can be expressed as a sum of terms  $H = \sum_a H_a$ . Then we may approximate the time evolution operator  $e^{-i\Delta H}$  for time interval  $\Delta$  using

$$\begin{aligned} & e^{-i\Delta(H_1+H_2+\dots H_n)} - \left( e^{-i\Delta H_1} e^{-i\Delta H_2} \dots e^{-i\Delta H_n} \right) \\ &= \frac{1}{2} \Delta^2 \sum_{a < b} [H_a, H_b] + \dots \end{aligned} \quad (21)$$

If the Hamiltonian  $H$  is geometrically local, then  $e^{-i\Delta H_a}$  acts on a constant number of qubits, and let us therefore assume it can be simulated accurately with a constant number of gates (ignoring a possible

“Solovay-Kitaev slowdown”). Furthermore, for an  $n$ -qubit system, the total number of terms  $\{H_a\}$  is  $O(n)$ , and the number of terms in  $\{H_a\}$  that fail to commute with any fixed term  $H_b$  is a constant. Therefore, if all the terms have a bounded operator norm,  $\|H_a\| \leq E$ , then the error on the right-hand side of eq. (21) has operator norm no larger than  $C\Delta^2 E^2 n$ , where  $C$  is a constant. To simulate the evolution for time  $T$ , we need  $T/\Delta$  time steps, and the accumulated error is bounded above by  $CT\Delta E^2 n$ . Therefore our simulation has constant accuracy if we choose  $\Delta = O(1/nT)$  (assuming  $E$  is a constant), and circuit size  $O(nT/\Delta) = O[(nT)^2]$ . In effect, then, the circuit size scales quadratically with the volume of the simulated spacetime.

The object of this exercise is to show that the cost of the simulation can be reduced to a lower power of  $nT$ , namely the  $3/2$  power.

a) Show that

$$e^{2(A+B)} - e^A e^B e^B e^A = \frac{1}{3}[A, [A, B]] - \frac{2}{3}[B, [B, A]] + \dots, \quad (22)$$

where now the ellipsis indicates terms of quartic or higher order.

b) Show that

$$\begin{aligned} & e^{-i\Delta(H_1+H_2+\dots+H_n)} \\ & - \left( e^{-i\Delta H_1/2} e^{-i\Delta H_2/2} \dots e^{-i\Delta H_n/2} \right) \\ & \times \left( e^{-i\Delta H_n/2} \dots e^{-i\Delta H_2/2} e^{-i\Delta H_1/2} \right) \\ & = \Delta^3 \left( \sum_{a<b<c} \mathcal{O}_{abc} + \sum_{a<b} \mathcal{O}_{ab} + \dots \right), \quad (23) \end{aligned}$$

where the ellipsis indicates terms of quartic order and above. Find explicit expressions for  $\mathcal{O}_{abc}$  and  $\mathcal{O}_{ab}$  in terms of third-order commutators of the terms in  $\{H_a\}$ .

c) Assuming as above that  $e^{-i\Delta H_a/2}$  can be simulated using a constant number of gates, show that for a geometrically local  $n$ -qubit Hamiltonian  $H = \sum_a H_a$ , where  $\|H_a\| \leq \text{constant}$ , the time evolution operator  $e^{-iHT}$  can be simulated to constant accuracy using  $O[(nT)^{3/2}]$  gates.

The power of  $nT$  can be reduced further using more elaborate constructions.